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Thermal uncertainty quantification in frequency responses of laminated composite plates

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Abstract

The propagation of thermal uncertainty in composite structures has significant computational challenges. This paper presents the thermal, ply-level and material uncertainty propagation in frequency responses of laminated composite plates by employing surrogate model which is capable of dealing with both correlated and uncorrelated input parameters. The present approach introduces the generalized high dimensional model representation (GHDMR) wherein diffeomorphic modulation under observable response preserving homotopy (D-MORPH) regression is utilized to ensure the hierarchical orthogonality of high dimensional model representation component functions. The stochastic range of thermal field includes elevated temperatures up to 375K and sub-zero temperatures up to cryogenic range of 125K. Statistical analysis of the first three natural frequencies is presented to illustrate the results and its performance.

Keywords: A. Laminates; B. Vibration; C. Computational modelling; C. Statistical properties/methods; Thermal uncertainty

1. Introduction

Composite materials are being increasingly utilized in aerospace applications due to high strength, stiffness, light weight and tailorable properties. They may be exposed to variation in environmental (hot or cold) conditions during the service life such as aircraft wing made of composite materials experiences a wide range of temperature variation from take-off to level-
The free vibration of laminated plates with effect of environment has been considered earlier by Whitney and Ashton [1]. The effect of environments on the material properties of composites was studied by many researchers, for example, Strife and Prewo [2] and Bowles and Tompkins [3] and Seng et al. [4]. Ample published work is found on deterministic buckling analysis in conjunction to thermal and hygrothermal behavior [5-12]. The concept of random vibration is exhaustively utilised in many engineering application [13-15]. Most of the literatures are deterministic in nature, which lacks in portraying the probable deviation caused by random input parameters. Due to presence of large number of inter-dependent factors in production of composites, the system input parameters are generally random in nature. The allowable responses for conventional material is expected to be close to their
mean values as fewer parameters are involved in the ir production process while in contrast for composites, a range of random fluctuation in system parameter may occur due to large number of system properties. Even after ensuring the effective quality control of production process in sensitive applications, the allowable responses are normally found to scatter widely with respect to the mean values. The knowledge of input variabilities and corresponding range of stochastic responses may serve to control the purposes of lightweight design, which is one of the important characteristics for composites. Therefore the efficient computational modelling and analysis is needed considering randomness in material properties and ply orientation angle including the effect of thermal uncertainty to ensure optimization, operational safety and reliability. Such issues can be addressed by employing probabilistic method, which quantifies the uncertainties in frequency responses.

The novelty of the present study includes the stochastic analysis of natural frequencies for laminated composite plates subjected to uncertain thermo-mechanical loading. A surrogate model is employed by using the generalized high dimensional model representation (GHDMR) approach wherein D-MORPH (Diffeomorphic Modulation under Observable Response Preserving Homotopy) regression is employed to ensure the hierarchical orthogonality of HDMR component functions [16]. Random sampling high dimensional model representation (RS-HDMR) was employed for uncertainty quantification of natural frequency in composite plates considering three input parameters namely fibre-orientation angle, elastic modulus and mass density [17], wherein the input parameters are independent to each other. In contrast, investigation is also carried out with a new element for laminated composite plates [18] while stochastic modeling of unidirectional composites is studied considering delamination [19]. In the present study, the sources of uncertainty for natural frequency are considered as layer wise variation of material properties, ply orientation angle and temperature. Material properties of fibre reinforced composites are temperature
Due to this reason, the input parameters become co-related to each other in case of layerwise combined variation of material properties, ply orientation angle and temperature. The co-related input parameters cannot be mapped for the corresponding output response using conventional high dimensional model representation (HDMR) approach (different variants of HDMR can be found in available literature such as Cut-HDMR [20], RS-HDMR [21], mp-Cut-HDMR [22], Multicut-HDMR [23], lp-RS-HDMR [24] depending primarily on sampling scheme [16]. The present GHDMR can efficiently take care of both independent as well as co-related input parameters under a relaxed vanishing condition. The extended bases are used as basis functions to approximated HDMR component functions and D-MORPH regression is used to determine the coefficients in the GHDMR algorithm. The application of GHDMR is the first attempt of its kind in realm of laminated composites to take into account the effect of both non-correlated and/or correlated input parameters. In the present study, a random variable approach is employed in conjunction to finite element formulation to figure out the random eigenvalue problem. The numerical results are shown for first three natural frequencies with individual and combined layerwise variation of the stochastic input parameters. The present probabilistic approach is validated with Monte Carlo simulation wherein a small random variation is considered as tolerance zone.

2. Governing equations

Consider a laminated composite cantilever plate as furnished in Figure 1(a,b) with thickness ‘t’ consisting of n number of thin lamina, the stress strain relations in the presence of temperature can be represented as [25]

\[
\begin{align*}
\{\sigma_x(\tilde{\omega})\} &= \begin{bmatrix}
Q_{11}(\tilde{\omega}) & Q_{12}(\tilde{\omega}) & Q_{16}(\tilde{\omega}) \\
Q_{12}(\tilde{\omega}) & Q_{22}(\tilde{\omega}) & Q_{26}(\tilde{\omega}) \\
Q_{16}(\tilde{\omega}) & Q_{26}(\tilde{\omega}) & Q_{66}(\tilde{\omega})
\end{bmatrix} \begin{bmatrix}
\epsilon_x^o - \epsilon_x^T(\tilde{\omega}) \\
\epsilon_y^o - \epsilon_y^T(\tilde{\omega}) \\
\gamma_{xy}^o - \gamma_{xy}^T(\tilde{\omega})
\end{bmatrix} \\
\{\sigma_y(\tilde{\omega})\} &= \begin{bmatrix}
Q_{11}(\tilde{\omega}) & Q_{12}(\tilde{\omega}) & Q_{16}(\tilde{\omega}) \\
Q_{12}(\tilde{\omega}) & Q_{22}(\tilde{\omega}) & Q_{26}(\tilde{\omega}) \\
Q_{16}(\tilde{\omega}) & Q_{26}(\tilde{\omega}) & Q_{66}(\tilde{\omega})
\end{bmatrix} \begin{bmatrix}
\epsilon_y^o - \epsilon_y^T(\tilde{\omega}) \\
\epsilon_x^o - \epsilon_x^T(\tilde{\omega}) \\
\gamma_{xy}^o - \gamma_{xy}^T(\tilde{\omega})
\end{bmatrix} \\
\tau_{xy}(\tilde{\omega}) &= \begin{bmatrix}
Q_{11}(\tilde{\omega}) & Q_{12}(\tilde{\omega}) & Q_{16}(\tilde{\omega}) \\
Q_{12}(\tilde{\omega}) & Q_{22}(\tilde{\omega}) & Q_{26}(\tilde{\omega}) \\
Q_{16}(\tilde{\omega}) & Q_{26}(\tilde{\omega}) & Q_{66}(\tilde{\omega})
\end{bmatrix} \begin{bmatrix}
\epsilon_x^o - \epsilon_x^T(\tilde{\omega}) \\
\epsilon_y^o - \epsilon_y^T(\tilde{\omega}) \\
\gamma_{xy}^o - \gamma_{xy}^T(\tilde{\omega})
\end{bmatrix}
\end{align*}
\] (1)

and
where, $\sigma_x$, $\sigma_y$, $\tau_{xy}$, $\tau_{xz}$, $\tau_{yz}$ are normal and shear stresses; $\varepsilon^o_x$, $\varepsilon^o_y$, $\gamma^o_{xy}$, $\gamma^o_{xz}$, $\gamma^o_{yz}$ are normal and shear strains. The $e^T_x$, $e^T_y$, $e^T_{xy}$ values are the thermal strain components due to temperature in $x$-$y$ reference axes, which are derived from the corresponding values in the fiber axes after applying the transformations expressed as

$$
\{ e^T (\tilde{\omega}) \} = \begin{bmatrix} e^T_x (\tilde{\omega}) \\ e^T_y (\tilde{\omega}) \\ e^T_{xy} (\tilde{\omega}) \end{bmatrix} = \begin{bmatrix} \cos^2 \theta (\tilde{\omega}) & \sin^2 \theta (\tilde{\omega}) \\ \sin^2 \theta (\tilde{\omega}) & \cos^2 \theta (\tilde{\omega}) \\ -2 \cos \theta (\tilde{\omega}) \sin \theta (\tilde{\omega}) & 2 \sin \theta (\tilde{\omega}) \cos \theta (\tilde{\omega}) \end{bmatrix} \begin{bmatrix} \alpha_1 \\ \alpha_2 \end{bmatrix} \Delta T (\tilde{\omega}) \tag{3}$$

where, $(\tilde{\omega})$ indicates the randomness of the corresponding variables and $\alpha_1$, $\alpha_2$ are the thermal expansion coefficients of lamina in longitudinal and lateral directions and their values are considered as $-0.3 \times 10^6$ /K and $28.1 \times 10^6$ /K, respectively. $\Delta T (\tilde{\omega}) = T(\tilde{\omega}) - T_o (\tilde{\omega})$, where $T_o (\tilde{\omega})$ is the reference variable temperature in Kelvin. $T$ is exposed random temperature in Kelvin. Here, $\theta(\tilde{\omega})$ denotes the random ply orientation angle of the lamina with reference to $x$-axis. The non-mechanical in-plane stress and moment resultants due to thermal environment are expressed as

$$
\begin{bmatrix} N^t_x (\tilde{\omega}) \\ N^t_y (\tilde{\omega}) \\ N^t_{xy} (\tilde{\omega}) \end{bmatrix} = \sum_{k=1}^{n} \int_{Z_{k-1}}^{Z_k} [\overline{Q}_{ij} (\tilde{\omega})]_k [e^T (\tilde{\omega})]_k \, dz \quad \text{where } i, j = 1, 2, 6 \tag{4}
$$

and

$$
\begin{bmatrix} M^t_x (\tilde{\omega}) \\ M^t_y (\tilde{\omega}) \\ M^t_{xy} (\tilde{\omega}) \end{bmatrix} = \sum_{k=1}^{n} \int_{Z_{k-1}}^{Z_k} [\overline{Q}_{ij} (\tilde{\omega})]_k [e^T (\tilde{\omega})]_k \, z \, dz \quad \text{where } i, j = 1, 2, 6 \tag{5}
$$
The force and moment resultants are modified to include the thermal field by the constitutive equations [26, 27] for the composite plate are given by

\[
\{ F(\omega) \} = [D(\omega)] \{ \varepsilon \} - \{ F^0(\omega) \}
\]  

(6)

where

\[
\{ F(\omega) \} = [N_x(\omega) \quad N_y(\omega) \quad N_{xy}(\omega) \quad M_x(\omega) \quad M_y(\omega) \quad M_{xy}(\omega) \quad Q_x(\omega) \quad Q_y(\omega)]^T
\]

(7)

\[
[D(\omega)] =
\begin{bmatrix}
A_{11}(\omega) & A_{12}(\omega) & A_{16}(\omega) & B_{11}(\omega) & B_{12}(\omega) & B_{16}(\omega) & 0 & 0 \\
A_{12}(\omega) & A_{22}(\omega) & A_{26}(\omega) & B_{12}(\omega) & B_{22}(\omega) & B_{26}(\omega) & 0 & 0 \\
A_{16}(\omega) & A_{26}(\omega) & A_{66}(\omega) & B_{16}(\omega) & B_{26}(\omega) & B_{66}(\omega) & 0 & 0 \\
B_{11}(\omega) & B_{12}(\omega) & B_{16}(\omega) & D_{11}(\omega) & D_{12}(\omega) & D_{16}(\omega) & 0 & 0 \\
B_{12}(\omega) & B_{22}(\omega) & B_{26}(\omega) & D_{12}(\omega) & D_{22}(\omega) & D_{26}(\omega) & 0 & 0 \\
B_{16}(\omega) & B_{26}(\omega) & B_{66}(\omega) & D_{16}(\omega) & D_{26}(\omega) & D_{66}(\omega) & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & S_{44}(\omega) & S_{45}(\omega) \\
0 & 0 & 0 & 0 & 0 & 0 & S_{45}(\omega) & S_{55}(\omega)
\end{bmatrix}
\]

(8)

The non-mechanical loads due to uncertain thermal condition

\[
\{ F^0(\omega) \} = [N^0_x(\omega) \quad N^0_y(\omega) \quad N^0_{xy}(\omega) \quad M^0_x(\omega) \quad M^0_y(\omega) \quad M^0_{xy}(\omega) \quad 0 \quad 0 ]^T
\]

(9)

\[
\{ \varepsilon \} = [\varepsilon_x^o \quad \varepsilon_y^o \quad \gamma_{xy}^o \quad k_x \quad k_y \quad k_{xy} \quad \gamma_{xz}^o \quad \gamma_{yz}^o ]^T
\]

(10)

The stiffness coefficients are defined as [28]

\[
[A_{ij}(\omega), B_{ij}(\omega), D_{ij}(\omega)] = \sum_{k=1}^{n} \int_{Z_{k-1}}^{Z_k} \overline{Q}_{ij}(\omega) \, dz \quad \text{where } i, j = 1, 2, 6
\]

(11)

and

\[
[S_{ij}(\omega)] = \alpha_{sfc} \sum_{k=1}^{n} \int_{Z_{k-1}}^{Z_k} \overline{Q}_{ij}(\omega) \, dz \quad \text{where } i, j = 4, 5
\]

(12)

where \( \alpha_{sfc} \) is the shear correction factor and is assumed as 5/6. \( \overline{Q}_{ij}(\omega) \) in above equation (11) and (12) is defined as

\[
\overline{Q}_{ij}(\omega) = [T_i(\omega)]^{-1} [Q_{ij}(\omega)] [T_i(\omega)]^{-T} \quad \text{for } i, j = 1, 2, 6
\]

(13)

\[
\overline{Q}_{ij}(\omega) = [T_i(\omega)]^{-1} [Q_{ij}(\omega)] [T_i(\omega)]^{-T} \quad \text{for } i, j = 4, 5
\]

(14)
\[
\begin{bmatrix}
-Cos\theta(\tilde{\omega})Sin\theta(\tilde{\omega}) & Sin\theta(\tilde{\omega})Cos\theta(\tilde{\omega}) & Cos^2\theta(\tilde{\omega}) - Sin^2\theta(\tilde{\omega})
\end{bmatrix}
\]

and 
\[
[T_2(\tilde{\omega})] = \begin{bmatrix}
Cos\theta(\tilde{\omega}) & -Sin\theta(\tilde{\omega}) \\
Sin\theta(\tilde{\omega}) & Cos\theta(\tilde{\omega})
\end{bmatrix}
\]

\[
[Q_{ij}(\tilde{\omega})]_k = \begin{bmatrix}
Q_{11}(\tilde{\omega}) & Q_{12}(\tilde{\omega}) & 0 \\
Q_{12}(\tilde{\omega}) & Q_{22}(\tilde{\omega}) & 0 \\
0 & 0 & Q_{66}(\tilde{\omega})
\end{bmatrix}
\]

and 
\[
[Q_{ij}(\tilde{\omega})]_k = \begin{bmatrix}
Q_{44}(\tilde{\omega}) & 0 \\
0 & Q_{55}(\tilde{\omega})
\end{bmatrix}
\]

where \(Q_{ij}\) are the in-plane element of the stiffness matrix. From Hamilton’s principle [29], the dynamic equilibrium equation (for free vibration) can be expressed as \([30, 31]\)

\[
[M_e(\tilde{\omega})][\ddot{\delta}_e] + [K_e(\tilde{\omega})] \{\delta_e\} = 0
\]

(15)

where \(M_e(\tilde{\omega})\) and \(K_e(\tilde{\omega})\) are element mass and stiffness matrices, respectively. Here \([K_e(\tilde{\omega})] = [K_{ela}(\tilde{\omega})] + [K_{geo}(\tilde{\omega})]\) as the sum of element elastic stiffness matrix \([K_{ela}(\tilde{\omega})]\) and geometric stiffness matrix \([K_{geo}(\tilde{\omega})]\). After assembling all the element matrices and the force vectors with respect to the common global coordinates, the resulting equilibrium equation is formulated. Considering randomness of input parameters like temperature, ply-orientation angle, elastic modulus etc., the equation of motion of free vibration system after of assembling with \(n\) degrees of freedom can expressed as

\[
[M(\tilde{\omega})][\ddot{\delta}] + [K(\tilde{\omega})] \{\delta\} = 0
\]

(16)

where \(\{\delta\}\) denotes the vector of generalized coordinates. The governing equations are derived based on Mindlin’s Theory \([32]\) incorporating rotary inertia, transverse shear deformation. The random natural frequencies \([\omega_n(\tilde{\omega})]\) can be calculated employing standard eigenvalue problem \([33]\) and by solving the QR iteration algorithm. In the present study, an eight noded
isoparametric quadratic element with five degrees of freedom at each node (three translations and two rotations) is considered for finite element formulation with respect to laminated composite cantilever plate wherein the shape functions ($N_i$) are as follows

\begin{align}
N_i &= (1 + \xi \xi_i) (1 + \zeta \zeta_i) (\xi \xi_i + \zeta \zeta_i - 1) / 4 \quad \text{(for } i = 1, 2, 3, 4) \quad (17) \\
N_i &= (1 - \zeta^2) (1 + \zeta \zeta_i) / 2 \quad \text{(for } i = 5, 7) \quad (18) \\
N_i &= (1 - \xi^2) (1 + \xi \xi_i) / 2 \quad \text{(for } i = 6, 8) \quad (19)
\end{align}

where $\zeta$ and $\xi$ are the local natural coordinates of the element. The element stiffness matrix is given by

\begin{equation}
[K_e(\tilde{\omega})] = \int \int [B]^T [D(\tilde{\omega})][B][J_c]d\xi d\zeta \quad (20)
\end{equation}

where, $[B]$ is the strain displacement matrix and $[D(\tilde{\omega})]$ is the random stress-strain matrix.

The strain displacement matrix, $[B] = [[B_1], [B_2], \ldots \ldots [B_8]]$

\[
[B_i] = \begin{bmatrix}
N_{i,x} & 0 & 0 & 0 & 0 \\
0 & N_{i,y} & 0 & 0 & 0 \\
N_{i,y} & N_{i,x} & 0 & 0 & 0 \\
0 & 0 & 0 & N_{i,x} & 0 \\
0 & 0 & 0 & 0 & N_{i,y} \\
N_{i,y} - N_{i,x} & 0 & N_{i,y} & N_{i,x} & 0 \\
0 & 0 & N_{i,x} & N_{i} & 0 \\
0 & 0 & N_{i,y} & 0 & N_{i}
\end{bmatrix} \quad (21)
\]

The element mass matrix is obtained from the Integral

\begin{equation}
[M_e(\tilde{\omega})] = \int \int [N]^T [\rho(\tilde{\omega})][N][J_c]d\xi d\zeta \quad (22)
\end{equation}

where, $[N]$ is the shape function matrix and $[\rho(\tilde{\omega})]$ is the random inertia matrix. The derivatives of the shape function, $N_i$ with respect to $x, y$ are expressed in term of their derivatives with respect to $\xi$ and $\zeta$ by the following relationship
3. Formulation of GHDMR

The general high dimensional model representation (GHDMR) is important because in real practical applications, the variables are often correlated, for example, the cases wherein the input variables have some relations between them. Here relation can be deterministic or stochastic. For instance, large values of certain input variables may imply large or small values of some other stochastic input variables. Such relation may be controlled by some known or unknown distributions. These correlations are implicitly contained in the collected samples in practice. The GHDMR can construct a proper model for prediction of the random output (say natural frequency) in the stochastic domain. The present approach can treat both independent and correlated input variables, and includes independent input variables as a special case. The role of D-MORPH is to ensure the component functions’ orthogonality in hierarchical manner. The present technique decomposes the function $\lambda(S)$ with component functions by input parameters, $S = (S_1, S_2, \ldots, S_{kk})$. As the input parameters are independent in nature, the component functions are specifically projected by vanishing condition. Hence, it has limitation for general formulation. In contrast, a novel numerical analysis with component functions is portrayed in the problem of present context wherein a unified framework for general HDMR dealing with both correlated and independent variables are established. For different input parameters, the output is calculated as \[\lambda(S) = \lambda_0 + \sum_{i=1}^{kk} \lambda_i(S_i) + \sum_{1 \leq i < j \leq kk} \lambda_{ij}(S_i, S_j) + \ldots + \lambda_{12\ldots kk}(S_1, S_2, \ldots, S_{kk}) \quad (24)\]

\[\hat{\lambda}(S) = \sum_{u \subseteq kk} \hat{\lambda}_u(S_u) \quad (25)\]
\( \lambda_{12\ldots k} (S_1, S_2, \ldots, S_{kk}) \) indicates the residual contribution by input parameters. The subset \( u \subseteq \{1, 2, \ldots, kk\} \) denotes the subset where \( u \subseteq kk \) for simplicity and empty set, \( \varGamma \in u \). As per Hooker’s definition, the correlated variables are expressed as,

\[
\{ \lambda_u (S_u | u \subseteq kk) \} = \text{Arg min}_{g_u \in L^2(\mathbb{R}^k, u \subseteq kk)} \int \left( \sum_{u \subseteq k} g_u (S_u) - \lambda (S) \right)^2 w(S) dS \quad (26)
\]

\[
\forall u \subseteq kk , \quad \forall i \in u , \quad \int \lambda_u (S_u) w(S) dS_i dS_{-u} = 0 \quad (27)
\]

and

\[
\forall v \subseteq u , \quad \forall g_v : \int \lambda_u (S_u) g_v (S_v) w(S) dS = \langle \lambda_u (S_u), g_v (S_v) \rangle = 0 \quad (28)
\]

The function \( \lambda (S) \) can be obtained from sample data by experiments or by modelling. To minimise the computational cost, the reduction of the squared error can be realised easily.

Assuming \( H \) in Hilbert space is expanded on the basis \( \{h_1, h_2, \ldots, h_{kk}\} \), the bigger subspace \( \bar{H} (\supset H) \) is expanded by extended basis \( \{h_1, h_2, \ldots, h_{kk}, h_{kk+1}, \ldots, h_m\} \). Then \( \bar{H} \) can be decomposed as

\[
\bar{H} = H \oplus H^\perp \quad (29)
\]

where \( H^\perp \) denotes the complement subspace (orthogonal) of \( H \) [34] within \( \bar{H} \). In the past work [35-37], the component functions are calculated from basis functions. The component functions of Second order HDMR expansion are estimated from basis functions \( \{ \phi \} \) as [21]

\[
\lambda_i (S_i) \approx \sum_{r=1}^{kk} \alpha_r^{(0)i} \phi_r^i (S_i) \quad (30)
\]

\[
\lambda_{ij} (S_i, S_j) \approx \sum_{r=1}^{kk} [\alpha_r^{(i)i} \phi_r^i (S_i) + \alpha_r^{(j)i} \phi_r^j (S_j)] + \sum_{p=1}^{l} \sum_{q=1}^{l} \beta_p^{(0)ij} \phi_p^i (S_i) \phi_q^j (S_j) \quad (31)
\]

\( \lambda_{ij} (S_i, S_j) \approx \sum_{r=1}^{kk} [\alpha_r^{(i)i} \phi_r^i (S_i) + \alpha_r^{(j)i} \phi_r^j (S_j)] + \sum_{p=1}^{l} \sum_{q=1}^{l} \beta_p^{(0)ij} \phi_p^i (S_i) \phi_q^j (S_j) \quad (31)\)
equation system

$$\Gamma J = \hat{R}$$

(32)

where $\Gamma$ denotes a matrix ($N_{samp} \times \tilde{t}$) whose elements are basis functions at the $N_{samp}$ values of $S$; $J$ is a vector with $\tilde{t}$ dimension of all unknown combination coefficients; $\hat{R}$ is a vector with $N_{samp}$-dimension wherein $l$-th element is $\lambda(S^{(l)}) - \lambda_0$. $S^{(l)}$ denotes the $l$-th sample of $S$, and $\lambda_0$ represents the average value of all $\lambda(S^{(l)})$. The regression equation for least squares of the above equation can be expressed as

$$\frac{1}{N_{samp}} \Gamma^T \Gamma J = \frac{1}{N_{samp}} \Gamma^T \hat{R}$$

(33)

Due to the use of extended bases, some rows of the above equation are identical and can be removed to give an underdetermined algebraic equation system

$$AJ = \hat{V}$$

(34)

It has many of solutions for $J$ composing a manifold $Y \in \mathbb{R}^{\tilde{t}}$. Now the task is to find a solution $J$ from $Y$ to force the HDMR component functions satisfying the hierarchical orthogonal condition. D-MORPH regression provides a solution to ensure additional condition of exploration path represented by differential equation

$$\frac{dJ(l)}{dl} = \chi \nu(l) = (I_t - A^+ A) \nu(l)$$

(35)

wherein $\chi$ denotes orthogonal projector ensuring

$$\chi^2 = \chi \quad \text{and} \quad \chi^T = \chi$$

(36)

$$\chi = \chi^2 = \chi^T \chi$$

(37)
The free function vector may be selected to ensure the wide domain for $J(l)$ as well as to simultaneously reduce the cost $\kappa$ which can be expressed as

$$v(l) = -\frac{\partial \kappa(J(l))}{\partial J}$$

Then we obtain

$$\frac{\partial \kappa(J(l))}{\partial l} = \left(\frac{\partial \kappa(J(l))}{\partial J}\right)^T \frac{\partial J(l)}{\partial l} = \left(\frac{\partial \kappa(J(l))}{\partial J}\right)^T P \mathbf{v}(l)$$

$$= -\left(P \frac{\partial \kappa(J(l))}{\partial J}\right)^T \left(P \frac{\partial \kappa(J(l))}{\partial J}\right) \leq 0$$

The cost function can be expressed in quadratic form as

$$\kappa = \frac{1}{2} J^T B J$$

where $B$ denotes the positive definite symmetric matrix and $J_\infty$ can be expressed as

$$J_\infty = V_t (U_{\tilde{t}-r}^T V_{\tilde{t}-r})^{-1} U_{\tilde{t}-r}^T A^r \hat{V}$$

where the last columns ($\tilde{t}-r$) of $U$ and $V$ are denoted as $U_{\tilde{t}-r}$ and $V_{\tilde{t}-r}$ which can found by decomposition of $\chi B$ [38]

$$\chi B = U \begin{bmatrix} S_r & 0 \\ 0 & 0 \end{bmatrix} V^r$$

This unique solution $J_\infty$ in $Y$ indicates the minimized cost function. D-MORPH regression is used to find the $J$ which ensures the HDMR component functions’ orthogonality in hierarchical manner. The construction of the corresponding cost function $\kappa$ can be found in previous literature [16].

4. Random input representation

The random input parameters such as ply-orientation angle and temperature in each layer of laminate are considered for composite cantilever plates. It is assumed that the distribution of
random input parameters exists within a certain band of tolerance with their crisp values. The cases wherein the input variables considered in each layer of laminate are as follows:

(a) Variation of ply-orientation angle only: \[ \theta_i(\bar{\omega}) = \{\theta_1, \theta_2, \ldots, \theta_i \} \]

(b) Variation of longitudinal elastic modulus only: \[ E_1(\bar{\omega}) = \{E_{1(i)} \ldots E_{1(i)} \} \]

(c) Variation of shear modulus only: \[ G_{12}(\bar{\omega}) = \{G_{12(i)} \ldots G_{12(i)} \} \]

(d) Variation of temperature only: \[ T(\bar{\omega}) = \{T_{(1)} \ldots T_{(i)} \} \]

(e) Combined variation of ply orientation angle, elastic modulus, shear modulus and temperature: \[ [\theta, E_1, G_{12}, T(\bar{\omega})] = [(\theta_1 \ldots \theta_i), (E_{1(i)} \ldots E_{1(i)}), (G_{12(i)} \ldots G_{12(i)}), (T_{(1)} \ldots T_{(i)})] \]

where \( \theta_i, E_{1(i)}, G_{12(i)} \) and \( T_{(i)} \) are the ply orientation angle, temperature, respectively and ‘\( l \)’ denotes the number of layer in the laminate. In present study, \( \pm 5^\circ \) variation for ply orientation angle, \( \pm 10\% \) volatility in material properties and \( \pm 25K \) tolerance for temperature, respectively are considered from their respective deterministic values. Figure 2 presents the flowchart of frequency responses using GHDMR with D-MORPH. It is worth mentioning that material properties such as \( E_1 \) and \( G_{12} \) are considered as temperature dependant in the present study. Thus in case of the combined variation of ply orientation angle, elastic modulus, shear modulus and temperature, correlated input variables are needed to be mapped for natural frequencies as discussed in section-1.

5. Results and Discussion

The present study considers four layered graphite-epoxy angle-ply \([(\theta^\circ/-\theta^\circ/\theta^\circ/-\theta^\circ)]\) and cross-ply \((0^\circ/90^\circ/0^\circ/90^\circ)\) composite cantilever plates. An eight noded isoparametric plate bending element is considered for finite element formulation. Due to paucity of space, only a few important representative results are furnished. Table 1 presents the convergence study of non-dimensional fundamental natural frequencies of three layered graphite-epoxy untwisted composite plates [39]. Table 2 presents the non-dimensional natural frequencies for simply-
supported graphite-epoxy symmetric cross-ply composite plates \[40\]. In both the cases, close agreement with benchmarking results are obtained at \((6 \times 6)\) mesh size. Material properties and their variation with temperature \[41\] are furnished in Table 3. Considering mean temperature \(T=300\text{K}\) and thickness \(t=0.004\text{ m}\), the deterministic values of material properties are considered as \(E_1=E_2=15.4\text{ GPa}, \nu=0.43, G_{12}=G_{13}=G_{23}=3.56\text{ GPa}, \rho=1660\text{ kg/m}^3\). The present GHDMR methodology is employed to find a predictive and representative surrogate model relating each natural frequency to a number of input variables. The present surrogate models are used to determine the first three natural frequencies corresponding to given values of input variables, instead of time-consuming deterministic finite element analysis. The probability density function is plotted as the benchmark of bottom line results. Due to paucity of space, only a few important representative results are furnished. The variation of temperature is scaled in the range with the lower and the upper limit (tolerance limit) as \(\pm25\text{K}\) with respective mean values while for ply orientation angle as within \(\pm5^\circ\) fluctuation (as per standard of composite manufacturing industry) with respective deterministic values. Both angle-ply and cross-ply composite cantilever plates are considered for the present analysis.

Table 1 Convergence study for non-dimensional fundamental natural frequencies \([\omega=\omega_n L^2/\sqrt{(\rho/E_1 t^2)}]\) of three layered \((\theta^\circ/-\theta^\circ/\theta^\circ)\) graphite-epoxy untwisted composite plates, \(a/b=1, b/t=100\), considering \(E_1 = 138\text{ GPa}, E_2 = 8.96\text{ GPa}, G_{12} = 7.1\text{ GPa}, \nu_{12} = 0.3\).

<table>
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<th>Ply angle, (\theta)</th>
<th>Present FEM ((4 \times 4))</th>
<th>Present FEM ((6 \times 6))</th>
<th>Present FEM ((8 \times 8))</th>
<th>Present FEM ((10 \times 10))</th>
<th>Qatu and Leissa [39]</th>
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<tr>
<td>0°</td>
<td>1.0112</td>
<td>1.0133</td>
<td>1.0107</td>
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<td>0.4553</td>
<td>0.4549</td>
<td>0.4613</td>
</tr>
<tr>
<td>90°</td>
<td>0.2553</td>
<td>0.2567</td>
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<td>0.2542</td>
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Table 2 Non-dimensional natural frequencies $[\omega = \omega_n a^2 \sqrt{\frac{\rho}{E^2 t^2}}]$ for simply-supported graphite-epoxy symmetric cross-ply (0°/90°/90°/0°) composite plates considering $a/b=1$, $T=325K$, $a/t=100$.

<table>
<thead>
<tr>
<th>Frequency</th>
<th>Present FEM $(4 \times 4)$</th>
<th>Present FEM $(6 \times 6)$</th>
<th>Present FEM $(8 \times 8)$</th>
<th>Present FEM $(10 \times 10)$</th>
<th>Sai Ram and Sinha [40]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>8.041</td>
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<td>18.684</td>
<td>18.552</td>
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<td>38.597</td>
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</table>

Table 3 Material properties of glass/epoxy lamina at different temperatures, $E_1=E_2$, $G_{12}=G_{13}=G_{23}$, mass density ($\rho$)=1660 Kg/m$^3$, $\nu=0.43$ [41]

<table>
<thead>
<tr>
<th>Material properties (GPa)</th>
<th>Temperature (K)</th>
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<tr>
<td>$E_1$</td>
<td>15.4</td>
</tr>
<tr>
<td>$G_{12}$</td>
<td>3.56</td>
</tr>
</tbody>
</table>

Figure 3 presents the scatter plot which establishes the accuracy of present model with respect to original finite element model corresponding to random fundamental natural frequencies for combined variation of temperature and ply orientation angle. Table 4 presents the convergence study of present method compared to direct Monte Carlo simulation (MCS) for first three natural frequencies due to individual variation of ply-orientation angle and temperature of angle-ply ($45^\circ/-45^\circ/45^\circ/-45^\circ$) composite cantilever plate while Table 5 represents the convergence study of the present method with direct MCS for first three natural frequencies due to combined variation of temperature, ply-orientation angle, elastic modulus and shear modulus of angle-ply ($45^\circ/-45^\circ/45^\circ/-45^\circ$) composite cantilever plate. Figure 4(a-i) presents the comparative probability density function plot with respect to first three natural frequencies due to individual and combined variation of stochastic input parameters of angle-ply ($45^\circ/-45^\circ/45^\circ/-45^\circ$) composite cantilever plate for both MCS as well as present method. In present analysis, a sample size of 64 is considered for layerwise
meet the convergence criteria. Here, although the same sampling size as in direct MCS (10,000 samples) is considered, the number of actual FE analysis is much less compared to original MCS and is equal to number representative sample required to construct the surrogate model. The surrogate model is formed on which the full sample size of direct MCS is conducted. Hence, the computational time and effort expressed in terms of FE calculation is reduced compared to full scale direct MCS. This provides an efficient affordable way for simulating the uncertainties in natural frequency.
Table 4 Convergence study of first three natural frequencies due to individual variation of ply-orientation angle and temperature of angle-ply composite cantilever plate considering $E_1=E_2=15.4$ GPa, $G_{12}=G_{13}=G_{23}=3.56$ GPa, $\rho=1660$ Kg/m$^3$, $t=0.004$ m, $\nu=0.43$, $T_{\text{mean}}=300$K

| Parameter | Values | $f_1$ | | | $f_2$ | | | $T(\tilde{\omega})$ | | | $\theta(\tilde{\omega})$
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<td>Present Method (Sample run)</td>
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Table 5: Convergence study of the present method with direct Monte Carlo simulation (MCS) for first three natural frequencies due to combined variation of temperature, ply-orientation angle, elastic modulus and shear modulus of angle-ply (45°/-45°/45°/-45°) composite cantilever plate

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A comparative study on variation of stochastic natural frequencies is carried out for angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) composite cantilever plate due to individual variation of elastic modulus, shear modulus as furnished in Figure 5(a-f). From Figure 5, it is observed that the mean fundamental natural frequency for angle-ply laminate is found to be slightly lower than that of the same for cross-ply laminate while a significant higher mean values are obtained at second and third modes for angle-ply compared to cross-ply. Considering only variation of temperature of angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) composite cantilever plate, the probability density function (PDF) with respect to first three natural frequencies are plotted in Figure 6(a-f) wherein it is found that as
the temperature increases. This can be attributed to the fact that due to rise in temperature influences the thermo-mechanical loading due to random variation leading to change in the system properties. In contrast, the reverse trend is identified for cross-ply laminated composite plates due to reduction effect of 0° and 90° on effective stiffness of the laminate. On the other hand, Figure 7(a-c) presents the ply level quantification of uncertainty in first three natural frequencies in terms of Probability density function for angle-ply [(θ°/-θ°/θ°/-θ°) where θ=Ply orientation angle] and cross-ply (0°/90°/0°/90°) composite cantilever plate. Due to random variation of ply orientation angle, the elastic stiffness of the laminated composite plate is found to be varied which in turn influence the frequency responses irrespective of laminate configuration. The effect of combined variation of input parameters is also carried out in addition to individual variation of inputs in conjunction to stochastic natural frequencies for composite laminated plates as furnished in Figure 8(a-c). The ply orientation angle, elastic modulus, shear modulus and temperature of angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) are considered as random input variables wherein the upper and lower bounds of volatility in natural frequencies are found to be wider than that of individual variation of inputs irrespective of laminate configuration. This corroborates with the fact that the combined effect of random input parameters leads to increase the variation in outputs compared to individual cases.

In the present study, the relative coefficient of variance (RCV) (normalized mean to standard deviation ratio) due to variation of temperature is also quantified for each layer for angle-ply and cross-ply laminate as furnished in Figure 9(a,b). The two outer-most layers of the angle-ply laminate is found to be most sensitive to temperature variation for first three modes while the maximum sensitiveness of temperature is observed only at bottom layer of
layerwise ply-level sensitiveness to temperature variation for fundamental mode is studied to map the sensitivity of each layer due to the influence of ply orientation angle on variation of temperature as furnished in Figure 10(a-d). The least sensitivity is observed at $\theta(\bar{\omega})=45^\circ$ for outer layers of the angle-ply laminate.

6. Conclusions

This present study illustrates the layer-wise thermal uncertainty propagation with laminated composite plates. The ranges of variation in first three natural frequencies are analyzed considering both individual and combined stochasticity of input parameters. A generalized high dimensional model representation (GHDMR) model in conjunction with diffeomorphic modulation under observable response preserving homotopy (D-MORPH) regression is employed to map the input parameters (both correlated and uncorrelated) and natural frequencies. After utilizing the aforementioned surrogate modelling approach, the number of finite element simulations is found to be exorbitently reduced compared to original Monte Carlo simulation without compromising the accuracy of results. The computational expense is reduced by (1/156) times (individual stochasticity) and (1/19) times (combined stochasticity) of original Monte Carlo simulation. It is observed that as the temperature increases the variabilities of first three natural frequencies of angle-ply laminate are increased and the probability density function become steeper. The two outer-most layers of the angle-ply laminate is found to be most sensitive to temperature variation for first three modes while the maximum sensitivity of temperature is observed at bottom layer of the cross-ply laminate. It is found that stochastic variation of temperature influences the natural frequencies and thus it is a crucial design parameter from the operational safety and serviceability point of view. The numerical results obtained in this study provide a comprehensive idea for design and
Acknowledgement

HR and GL acknowledge the support of ONR, USA with account number N00014-11-1-0716. SA acknowledges the support of Royal Society of London through the award of Wolfson Research Merit award.

References


Fig. 1(a) Laminated composite cantilever plate

Fig. 1(b) Force and moment resultants diagram
Identification and definition of range of input parameters
\(T(\bar{\omega}), \theta(\bar{\omega}), E_{1}(\bar{\omega}), G_{12}(\bar{\omega})\)

Finite Element formulation to evaluate natural frequencies

Selection of design points based on Sobol sequence Sampling

Mapping frequency responses using FEM

Model construction using GHDMR algorithm

Model verification by PDF, Convergence study

MCS using GHDMR

Statistical analysis for natural frequency

Fig. 2 Flowchart of frequency responses using GHDMR with D-MORPH
Fig. 3 Scatter plot for present surrogate model with respect to original FE model of fundamental natural frequencies for combined variation of ply-orientation angle and temperature of graphite-epoxy angle-ply (45°/-45°/45°/-45°) composite cantilever plate, considering $E_1 = E_2 = 15.4$ GPa, $G_{12} = G_{13} = G_{23} = 3.56$ GPa, $T = 300$ K, $\rho = 1660$ Kg/m$^3$, $t = 0.004$ m, $\nu = 0.43$. 
Parameters

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<th>Parameters</th>
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<th>$f_2$</th>
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<td><img src="image" alt="Graph of $f_2$ for $\theta(\tilde{\omega})$" /></td>
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<td><img src="image" alt="Graph of $f_2$ for $T(\tilde{\omega})$" /></td>
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<td>$[\theta,T,E_1,G_{12}(\tilde{\omega})]$ (Combined)</td>
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Fig. 4 (a-i) Probability density function with respect to first three natural frequencies (Hz) due to individual and combined variation of ply orientation angle, elastic modulus, shear modulus and temperature of angle-ply ($45^\circ/-45^\circ/45^\circ/-45^\circ$) composite cantilever plate at mean temperature ($T_{\text{mean}}$) = 300K.
Fig. 5 (a-f) Probability density function with respect to first three natural frequencies (Hz) due to individual variation of elastic modulus, shear modulus of angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) composite cantilever plate at mean temperature ($T_{\text{mean}}$) = 300K.
Fig. 6 (a-f) Probability density function with respect to first three natural frequencies (Hz) due to individual variation of temperature of angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) composite cantilever plate.
Fig. 7(a-c) Probability density function with respect to first three natural frequencies (Hz) due to individual variation of ply orientation angle of angle-ply \((\theta^\circ/-\theta^\circ/\theta^\circ/-\theta^\circ)\) where \(\theta=0^\circ, 15^\circ, 30^\circ, 45^\circ, 60^\circ, 75^\circ\) and \(90^\circ\) and cross-ply\((0^\circ/90^\circ/0^\circ/90^\circ)\) composite cantilever plate at mean temperature \(T_{\text{mean}}=300\text{K}\).
Fig. 8(a-c) Probability density function with respect to first three natural frequencies (Hz) due to combined variation of ply orientation angle, elastic modulus, shear modulus and temperature of angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) composite cantilever plate at mean temperature ($T_{\text{mean}}$) = 300K.
Fig. 9(a,b) Relative coefficient of variance (RCV) of first three natural frequencies due to variation of temperature (layerwise) for angle-ply (45°/-45°/45°/-45°) and cross-ply (0°/90°/0°/90°) composite cantilever plate at mean temperature ($T_{\text{mean}}$)=300K.

Fig. 10(a-d) Relative coefficient of variance (RCV) of fundamental mode due to variation of temperature (layerwise) for angle-ply ($\theta^\circ$) composite cantilever plate at mean temperature ($T_{\text{mean}}$)=300K.