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Dark electrical bias effect on moisture-induced degradation in inverted lead halide perovskite solar cells measured by advanced chemical probes

Jérémy Barbé†, Vikas Kumar‡, Michael J. Newman†, Harrison K.H. Lee§, Sagar M. Jain†, Hu Chen§, Cécile Charbonneau†, Cornelia Rodenburg‡, Wing C. Tsoi†*

†SPECIFIC, College of Engineering, Swansea University, Bay Campus, Fabian Way, Swansea, SA1 8EN, U.K.
‡Department of Materials Science and Engineering, University of Sheffield, Mappin Street, Sheffield S1 3JD, U.K.
§KAUST Solar Center, Physical Science and Engineering Division, King Abdullah University of Science and Technology, Thuwal 23955-6900, Saudi Arabia

AUTHOR INFORMATION

Corresponding Author

*E-mail: W.C.Tsoi@Swansea.ac.uk

†Electronic supplementary information (ESI) available.
ABSTRACT: Emerging lead halide perovskite materials have enormous potential for a range of optoelectronic devices, such as solar cells, light emitting diodes, transistors and lasers. However, the large-scale commercialization of these technologies will depend on the ability of the active material to be stable under environmental and operating conditions. In this work, we measured the first time the electrical bias-induced degradation of inverted perovskite solar cells in the dark in different environments and concluded that humidity coupled with electrical bias results in fast degradation of CH$_3$NH$_3$PbI$_3$ into PbI$_2$. Micro-Raman and photoluminescence show that the degradation starts from the edge of the cell due to moisture ingress. By using novel local Raman-transient photocurrent measurements, we were able to probe local ion migration at the degraded region and non-degraded region and found that the formation of PbI$_2$ can passivate perovskite by reducing ion migration. The degradation is far from uniform across different grains as revealed by secondary electron hyperspectral imaging, an advanced scanning electron microscopy technique which allows probing the composition of individual grain from the cross-section. By using potential step chronoamperometry, we also found that the bias degradation is closely related to the density of mobile ions. The unique combination of established methods with several novel analytical tools provides an insight into the origin of the bias-degradation of inverted perovskite solar cells from nano-scale to cell level, and demonstrates the potential of these novel tools for studying the degradation in other perovskite systems.

Introduction

Since their invention in 2009\textsuperscript{1}, organometal halide perovskite solar cells (PSCs) have shown unprecedented performance improvements, now exceeding 22% power conversion efficiency\textsuperscript{2}. Beyond their performance, perovskites are the subject of intense research interest due to their
potential for easier and lower-cost manufacturing than the market-leading silicon solar technology. However, ensuring the long-term stability of perovskite solar cells remains a challenge for commercialization of this technology. Perovskite solar cells tend to degrade prematurely depending on the films quality and atmospheric conditions. Humidity, light, oxygen, temperature and more recently electrical stress have been shown to play an important role in initiating degradation pathways. Among these, electrical stress in “dark” has been very rarely investigated despite being a critical parameter for solar cells. Besides, perovskites are also being considered for other optoelectronic devices such as field-effect light emitting diodes (LEDs), transistors or lasers, which also operate under an external bias. In particular, LEDs have an architecture and films thickness similar to solar devices but a relatively high turn-on voltage of 3-4 V. This implies that perovskites should be stable under high electric field to achieve maximum external quantum efficiency and luminance.

Various mechanisms have been proposed to explain the degradation of the perovskite phase and resulting device performance upon applying an electric field in the dark. Using a laterally contacted device, it was suggested that methylammonium cations are loosely bound in the hydrated perovskite phase formed in the presence of moisture. These cations drift and accumulate near the negative electrode under the effect of an electric field, which leads to the irreversible structural modification of the depleted perovskite into PbI$_2$ near the positive electrode. Ions migration was also observed by temperature-dependent dark current decay in a standard (top-anode) perovskite device architecture. The device stability was found to strongly depend on the level of applied voltage in the dark. Irreversible degradation was observed when $V_{\text{applied}} > V_{\text{oc}}$. This was explained by the inversion of the space charge region at the electron transport layer (ETL) interface after the displacement of a high density of ionic defects (ions or
ion vacancies). Unlike these observations, D. Bryant and co-workers showed that the rapid electrical bias-induced degradation of CH$_3$NH$_3$PbI$_3$ in a standard device architecture can be attributed to the diffusion of oxygen in the perovskite lattice and formation of highly reactive superoxide species O$_2^-$ when electrons are not extracted efficiently from the active layer$^{16}$. In this case, slower degradation was observed when devices were exposed to moisture alone. Hence a crucial question is to ascertain which mechanism is responsible for the degradation of PSCs under applied bias in “dark” and to clearly establish the role played by superoxides or/and ionic defects migration.

In this work, we prepared inverted PSCs using NiO$_x$ as HTL and investigated the effect of electrical bias in “dark” on the device performances and material properties. NiO$_x$ hole transport layer was chosen in this study because of its higher efficiency and reproducibility over more commonly used PEDOT:PSS. Moreover, in the course of this study, we also investigated typical PEDOT:PSS and organic small molecules as alternative HTLs and observed similar degradation, which tend to show that the degradation can be quite general. To the best of our knowledge, study on the bias degradation of inverted PSCs in dark is “missing” although they have several advantages over standard top-anode devices such as less hysteresis and lower processing temperatures, while record efficiencies are getting comparable to the standard architecture. Besides, important questions remain for the dark bias degradation such as: What is the chemical origin of such degradation? How does the device degrade structurally? Here, we provide some insights to these questions together with new and advanced structural probes.

Methods

1.1 Materials
All materials were used without purification. Anhydrous dimethyl sulfoxide (DMSO), N,N-dimethylformamide (DMF), anhydrous ethanol, anhydrous chlorobenzene, anhydrous ethanolamine, anhydrous 2-methoxyethanol, bathocuproine (BCP, 96%), nickel acetate (Ni-acc, 98%) and lead iodide (PbI₂, 99%) were purchased from Sigma-Aldrich. Methylammonium iodide (MAI) was purchased from Dyesol. The electron transport material [6,6]-phenyl-C₆₁ butyric acid methyl ester (PCBM) was obtained from Solenne BV. To prepare the NiOₓ solution, 0.2 M solution of nickel acetate tetrahydrate was dissolved in a 1:0.012 volume ratio of 2-methoxyethanol:ethanolamine and stirred for 1 h at 60°C. CH₃NH₃PbI₃ precursor solution was prepared by mixing CH₃NH₃I and PbI₂ (1:1.05 molar ratio) in DMF/DMSO (4:1 volume ratio) with a concentration of 804 mg mL⁻¹ at 60°C for 2 h. PCBM was dissolved in chlorobenzene with concentration 20 mg mL⁻¹ and stirred overnight at 60°C. BCP was dissolved in anhydrous ethanol with concentration 0.5 mg mL⁻¹. All solutions were filtered with 0.45 µm PTFE syringe filters.

1.2 Perovskite solar cell fabrication and testing

15 Ω.□⁻¹ ITO/glass substrates (Lumtec) were sequentially cleaned with detergent in DI water, acetone, and isopropyl alcohol in an ultrasonic bath. Subsequently, they were dried with N₂ and exposed to O₂ plasma for 10 min. 20 nm-thick NiOₓ layer was deposited via spin-coating at 4000 rpm for 30 s followed by 30 min sintering at 250°C. Subsequent layers were prepared under nitrogen in a glove box. The perovskite precursor solution was spin coated at 4000 rpm for 30 s. Chlorobenzene was drop-cast onto the rotating sample after 15 s of spinning to promote crystallization. The 400 nm perovskite film was annealed at 100°C for 10 min. After cooling, the PCBM solution was spin-coated onto the CH₃NH₃PbI₃ layer at 4000 rpm for 30 s, followed by spin-coating the BCP solution at 6000 rpm for 20 s. The thickness of the PCBM is 60 nm. The
thickness of the BCP film is too thin to be measured by profilometry or SEM. Finally, 100 nm thick silver counter electrodes were evaporated at $10^{-4}$ Torr using an Edwards 306 thermal evaporator, forming devices with an active area of 0.15 cm².

Current density–voltage (J–V) characterizations were performed by a Keithley 2400 source-meter unit under 0.8 sun illumination using a Newport 92193A-1000 solar simulator. Current-voltage sweeps were performed from both $V_{OC}$-to-$J_{SC}$ and vice versa at a rate of 0.1 V s⁻¹.

1.3 Material characterizations

The Raman and photoluminescence measurements were performed with a Renishaw Invia Raman system in backscattering configuration. A 532 nm laser and 50x objective were used (NA: 0.50, spot size $\approx$ 1 µm). For single point Raman measurements, a laser power of 0.6 mW and acquisition time of 30 s were used; For the Raman and PL line mapping, 0.6 mW/5 s and 0.06 µW/1 s laser power/acquisition time were used for each measurement point, respectively. 1800 gr/mm and 300 gr/mm gratings were used for Raman and photoluminescence, respectively. Local photocurrent measurements were acquired using the laser beam and X-Y scanning stage from the Raman system (Renishaw Invia, 532 nm, 50x objective) and a source-meter unit Keithley 236. The devices were contacted using a Linkam atmospheric chamber. The solar cell was held at short-circuit in the dark until the residual current dropped below 5 nA (which can take up to 60 s). Then, the laser shutter was open and the photocurrent was monitored over time. A low laser power of 0.3 mW was used to avoid any further degradation of the perovskite film by the laser beam during the measurement.
XRD measurements were carried out using a Bruker D8 Discover instrument with a CuKα beam (wavelength is 0.15418 nm) at 40 kV and 40 mA, scan parameters of 1 s/step at 0.03° of 2θ step size.

The basic principles of the experimental method to collect hyperspectral images (SEHI) have been reported by P. Kazemian et al. In this work, we used FEI Helios NanoLab G3 UC SEM. All the samples are imaged using through-lens detector (TLD) at a working distance of 4.0 mm with a beam current of 6.3 pA and an accelerating voltage of 1 kV. The required hyperspectral window was selected (0-6.0 eV) by setting the mirror electrode to 0 V. For all SEHI images, the dwell time was adjusted to 50 ns.

Chronoamperometry measurements were carried out using a Keithley 236 source-meter unit controlled by software. The dark current through the cell was continuously recorded during the applied voltage sequence: the cell was first left at short-circuit for 30 s to allow equilibrium to be reached. Then, a forward bias of 1 V was applied for 30 s before the device was switched to short-circuit mode (0 V) for 30 s (until the current stabilized to a few nanoamperes). The device was switched on and off successively several times. The same procedure was also carried out at 0.5 V and 0.85 V.

The total charge density was estimated from the current transients at short-circuit condition by integrating the current decay as described below:

\[
Q = \frac{1}{A \times d} \int \frac{I(t)}{e} \, dt \tag{1}
\]

Where Q is the charge density, A the surface of the electrode, d the thickness of the perovskite layer, e the electronic charge and I the current at short circuit conditions.
Results and Discussion

We first investigated the effect of electrical bias in dark coupled with water vapor and/or oxygen on the PSCs performance. PSCs with structure glass/ITO/NiO$_x$/CH$_3$NH$_3$PbI$_3$(MAPI)/PC$_{60}$BM/BCP/Ag were fabricated and subjected to 1 V forward bias in air (40% RH), dry N$_2$, dry air and humid N$_2$. All degradations in this work were performed in the dark. In the case of humid N$_2$, a water-saturated atmosphere (≈100% humidity) was obtained by flowing N$_2$ through a sealed water jar. The current-voltage (J-V) curves before and after bias applications are shown in Figure 1a. When exposed to air or humid N$_2$, the cell drastically degraded after 1 h bias application mostly due to a large drop in $J_{sc}$. When exposed to dry N$_2$ or dry air, the biased cells only suffered minor degradation characterized by a slight drop of the $J_{sc}$, but no change in FF or $V_{oc}$. These results demonstrate the key role played by water vapor in the degradation of PSCs under bias in the dark, whereas oxygen does not have a significant effect. Even in scenarios where devices were degraded in dry N$_2$ or dry air, the slight loss in efficiency (from 10.4% to 8.8% and 10.7% to 8.8%, respectively) may be attributed to the presence of small amounts of water (or any polar molecules as explained by Leijtens et al.$^{10}$) in the perovskite stack, due to the sample being exposed to the ambient atmosphere prior to conducting the degradation tests, or perhaps while in contact with polar solvents in the glovebox during device fabrication. Oxygen and the formation of superoxides, although suggested by others$^{16}$ as another important path towards the degradation of lead halide perovskite materials, doesn’t appear to have a significant effect in this case (dark, dry environment, 1 V). It was shown that oxygen molecules enter the perovskite at a fast rate and uniformly saturate perovskite vacancies after a few seconds only$^{16}$. As we purged the atmospheric chamber for 10 min before measurement with a relatively high flow rate of dry air, it is reasonable to assume that samples
exposed to this environment were saturated with oxygen when the bias was applied. The fact that
the device was relatively stable under dry air (same trend as in dry N₂) diminishes the role of
highly reactive superoxide species, at least during the first hour of degradation.

Figure 1. (a) Current density – voltage curves of perovskite devices at 0.8 sun before (solid blue
lines) and after (dashed blue lines) biasing at 1 V for 1 h under dark condition in different
atmospheres. Reference cells left in the same environment for the same duration but without bias
are also shown for comparison (black lines). (b) Corresponding Raman spectra for the biased
cells and reference PbI₂ film (intensity divided by 10 times for comparison).

Figure 1b shows the Raman spectra for the four different cells after degradation. The cells
degraded in 40% RH air and humid N₂ reveal the features of PbI₂, with a main peak at 97 cm⁻¹
associated to the stretching mode of Pb-I⁹ (the higher intensity measured for 40% RH is due to
local inhomogeneities but does not reflect the degradation level), and other characteristic peaks
at 75 cm⁻¹, 110 cm⁻¹, 167 cm⁻¹ and 215 cm⁻¹. This indicates that a large fraction of the MAPI
converted to PbI₂ when exposed to 40% RH air or humid N₂ under 1 V bias. In the case of
devices exposed to dry air or dry N\textsubscript{2}, the Raman data confirm very small amount of PbI\textsubscript{2}, thus correlating device performance results.

Numerous reports showed that metal contacts play an important role in the degradation of perovskite devices. In particular, it has been observed that in presence of moisture, Ag can be converted to AgI due to the migration of I\textsuperscript{-} ions and Ag through pinholes in the Spiro-OMeTAD layer in the standard architecture\textsuperscript{20}. However, we didn’t observe any XRD diffraction peaks due to AgI (peaks at $2\theta = 22.3^\circ$, 23.8$^\circ$, and 39.2$^\circ$ for $\beta$-AgI)\textsuperscript{21} when measuring the bias degraded cell (1 V, 1 h, humid N\textsubscript{2}) through the Ag top electrode (Figure S1a, blue curve). It suggests that such degradation mechanism is not likely to happen in the inverted perovskite device structure, and can be explained by the fundamental difference that unlike standard device structure, a positive bias is applied to the ITO and not to the Ag electrode for the inverted device structure (so negative ions will be attracted to the ITO). The only difference between the XRD spectra of the degraded and non-degraded cell is the appearance of PbI\textsubscript{2} peak at 12.6$^\circ$ after degradation, which agrees well with the Raman results.
Figure 2. (a) Optical microscope image of the “long” edge of the degraded perovskite cell after 20 min under 1 V bias in humid air. (b) PbI$_2$ Raman intensity profiles measured at 96 cm$^{-1}$ along the edge. (c) Photoluminescence peak intensity profile at 775 nm. (d) Optical image showing the ITO and Ag electrodes delimitating the cells measured in this work. The “long” and “short” edges of the cells are shown as red and blue lines, respectively.

To confirm the role of moisture in the degradation of PSCs, freshly prepared samples were exposed to humid N$_2$ environment for 20 min without bias. This procedure allowed water molecules to penetrate within the perovskite film, in particular where the stack is not covered by the Ag top electrode or near the edges of the electrode. After this preconditioning, a bias was applied in air for short times, and optical microscopic images of the cell edge were recorded at different intervals (Figure S2). No visual degradation is observed up to 70 s at 1 V bias, before exposure to humid N$_2$. However, following exposure to humidity, the application of a bias clearly triggers the appearance of degradation features after a few seconds only. These are particularly concentrated near the long edge of the cell while the center of the cell and short edges are much better preserved. In this paper, the “long” edges are defined by the silver top electrode, while the “short” edges are defined by the ITO stripe, as shown in Figure 2(d). Similar patterns were observed when the cell is exposed to air with 40% RH and biased for 20 min, as shown in Figure 2a. These observations were further correlated with Raman and photoluminescence intensity profiles acquired across 300 µm from outside the edge of a cell toward the center, as shown in Figure 2b and c. The results show that the cell is highly degraded across 130 µm, starting from the edge of the cell. This is supported by 3-fold increase of PbI$_2$ Raman peak at 96 cm$^{-1}$ and more than 10 times drop in PL intensity at 775 nm.
The degradation pattern could be explained as follows: the 100 nm-thick silver top electrode acts as a barrier for water vapor and protects, to some extent, the underlying perovskite film. However, moisture ingress is much stronger at the long edges of the cell where the silver electrode finishes. As water penetrates in the device stack from the long edge of the cell, hydrated phases of perovskite are formed (as shown by several reports\(^3,4,22\)), which then rapidly converts to PbI\(_2\) under the effect of the electric field (Figure S1). In comparison, when the cell was not previously hydrated, the bias-induced degradation happens in minutes instead of seconds because of the much lower density of water molecules presents in the MAPI film (Figure 2). We also observed that the short edges were much less degraded than the long edges. Indeed, the short edges are delimited by the ITO bottom electrode while the Ag electrode continues further away, thus better protecting the short edges from environmental conditions.

XRD measurements were also carried out to investigate the formation of hydrated phases of perovskite when exposed to a humid environment. For this purpose, a sample was measured by focusing the X-ray beam outside the Ag electrode, to directly probe the perovskite/PCBM/BCP film after exposure to high humidity environment (1 h, 100% RH in N\(_2\)). For repeated measurements on different samples, most of the time no diffraction peaks due to monohydrate CH\(_3\)NH\(_3\)PbI\(_3\):H\(_2\)O or dihydrate (CH\(_3\)NH\(_3\))\(_4\)PbI\(_6\):2H\(_2\)O were observed.\(^3\) However, these metastable intermediate phases are usually difficult to measure because they are time and space-limited. Even recent in-situ synchrotron XRD measurements in high humidity environment could detect hydrated phases only from time to time during the degradation process.\(^22\) Thus, the fact that we couldn’t easily detect such phase doesn’t mean that monohydrate or dihydrate perovskite is not formed when moisture penetrates within the film, especially as the samples were measured at least 15 min after high humidity exposure and could have dried in the meantime depending on
the atmospheric conditions (Leguy et al. showed that hydrated perovskite crystals can dry in less than 15 min when exposed to dry air\(^3\)). Note that, occasionally, we observed a new peak at 37.9°, which has similar position to the XRD peak of \((\text{CH}_3\text{NH}_3)_4\text{PbI}_6\cdot2\text{H}_2\text{O}\) (Figure S1b, orange curve\(^3\)). It tends to suggest that hydrated perovskite phases form in the film, although they are momentary and dependent on the atmospheric conditions.

**Figure 3.** (a) Optical microscope image of the degraded cell with the 2 different regions (edge and inner side of the cell) used for local characterizations (b –d). (b) Micro-Raman spectra. (c) Local photocurrent stabilization curves measured at short-circuit (laser excitation power ≈0.3 mW). (d) Normalized photocurrent stabilization curves.
At this stage, the results demonstrate that the combined effects of water molecules ingress and applied bias are responsible for the degradation of the perovskite film to PbI$_2$. Next, making use of the various degradation levels that are formed at the cell edge after biasing, we performed localized photoelectrical characterizations by focusing the laser beam of the Raman system (\(\leq 1 \mu m\) diameter) at the center (region 1) or at the edge (region 2) of the cell, as represented in Figure 3a. The Raman spectra shown in Figure 3b confirm that two distinct levels of degradation are obtained: for region 1, the CH$_3$NH$_3$PbI$_3$ phase was well preserved, showing two broad peaks at 250 cm$^{-1}$ and 110 cm$^{-1}$, which can be assigned to the methylammonium cations torsional and libration modes, respectively.$^{23}$ For region 2, significantly more PbI$_2$ (at 96 cm$^{-1}$ and 110 cm$^{-1}$) relative to CH$_3$NH$_3$PbI$_3$ is generated. Although a clear Raman signal from CH$_3$NH$_3$PbI$_3$ at 250 cm$^{-1}$ was not observed for region 2, it does not necessary mean that all perovskite degraded to form PbI$_2$, as the Raman cross-section of perovskite is much lower than that of PbI$_2$ and thus can be masked by the much stronger PbI$_2$ signal.$^{19}$ This is further supported by the significant amount of photocurrent generated at region 2 although PbI$_2$ absorbs very little (if any) at the laser excitation wavelength, as explained in the next paragraph.

Photocurrent stabilization curves for the two different regions are shown in Figure 3c and d. To ensure that all measurements were started from the same equilibrium condition, the solar cell was held at short-circuit in the dark until the residual current dropped below 5 nA (which can take up to 60 s). Then, the laser shutter was opened and the photocurrent was monitored over time. A low laser power of 0.3 mW was used to avoid any further degradation of the perovskite film by the laser beam during the measurement.$^{19}$ On the non-degraded region (region 1), the photocurrent quickly switches to -2.7 $\mu$A but then gradually rises to reach -4.8 $\mu$A after 80 s. On the contrary, for the highly degraded region (region 2), the photocurrent quickly reaches a
stabilized value of -1.4 µA without long transients. The long photocurrent transient can be explained by ion migration according to several recent works\textsuperscript{24,9,25}. The photovoltage induces mobile ions to slowly move in the perovskite layer and to screen the electric field until ions reach a new equilibrium. This slow motion of mobile ions (and vacancies) is observed by the rise of the photocurrent toward a steady-state value at short-circuit on turning on the illumination. Alternatively, Shao et al. explained the slow photocurrent transient by trap filling process at grain boundaries and perovskite surface\textsuperscript{26}.

The results shown in Figure 3(b) and (d) indicate that when a large amount of PbI\textsubscript{2} is formed, the photocurrent transient is suppressed: the photocurrent quickly switches to a stabilized value, similarly to crystalline silicon solar cells which don’t have ionic conductivity\textsuperscript{9}. The two different rising times for the photocurrent measured on regions 1 and 2 are more clearly observed on the normalized curves in Figure 3d. Since PbI\textsubscript{2} has very low absorption at 532 nm\textsuperscript{19}, it is likely that the photocurrent measured on region 2 is generated in residual perovskite which was masked in the Raman spectra by the large amount of PbI\textsubscript{2} formed. These results are in good agreement with recent reports showing that incorporating some excess PbI\textsubscript{2} into CH\textsubscript{3}NH\textsubscript{3}PbI\textsubscript{3} solution\textsuperscript{27} or forming PbI\textsubscript{2} secondary phase in CH\textsubscript{3}NH\textsubscript{3}PbI\textsubscript{3} film through post-deposition thermal annealing\textsuperscript{25} can significantly reduce J-V hysteresis by suppressing ion migration. Obviously, the amount of PbI\textsubscript{2} formed at the edge of the cell is much too high to improve the device performance in our case but the results indicate that ion migration can be suppressed or highly reduced in the residual perovskite as a large amount of PbI\textsubscript{2} is formed. Hence, the local Raman-photoelectrical characterizations of perovskite solar cells give direct physical insights on the correlation between the local degradation (amount of PbI\textsubscript{2}) and local electronic-ionic processes with micrometer
resolution. To our knowledge, the information obtained by this novel technique cannot be obtained by other approaches.

Figure 4. Cross-sectional hyperspectral secondary electron images (SEHI) in the range 0-6.0 eV before (a) and after (b) bias-induced degradation in 40% RH air (1 V for 1 h). (c) and (d) SEHI of the same devices and areas after a second scan.

Next, to better understand the effect of applied bias on the degradation of the perovskite layer in the presence of moisture, secondary electron hyperspectral cross-sectional images were acquired on a cell biased at 1 V in 40% RH air for 1 h, and compared to those of a reference cell (no bias).

Secondary electron hyperspectral imaging (SEHI) is an advanced imaging technique making use of differences in the secondary electrons spectra of materials in a SEM to characterize local variations in the composition and electronic structure of materials with sub-10 nm resolution. The hyperspectral images in Figure 4 show uniform contrast across the reference sample (Figure 4a), but clear contrast variations in the case of the degraded sample: some grains appear brighter than others (Figure 4b). It suggests a modification of the chemical composition of these grains, consistent with the hypothesis of ion migration. Interestingly, the contrast variation indicates that the degradation is not uniform at the sub-micrometer scale and varies when comparing individual
grains. We suggest that the degradation ability of each crystalline grain under bias depends on
the initial structural quality and composition of the grain. Indeed, the concentration of mobile
defects in individual grains can be affected by the processing conditions which often lead to local
non-uniformity. Here, based on the clear contrast variations observed between crystalline
grains from the same perovskite film, we may conclude that the application of an electrical bias
is responsible for the drift of variable concentrations of ionic defects within each grain, resulting
in distinct levels of highly localized degradation. At this stage of development of the technique
(SEHI), it is not possible to measure the precise composition of the degraded grains. However,
the Raman (and XRD) results suggest that these grains are composed of a large amount of PbI$_2$.
The fact that both degraded and non-degraded grains sit next to each other at the microscopic
scale in the highly degraded perovskite confirms that residual non-degraded perovskite remains
in the degraded cell, which correlates well with the local photoelectrical characterizations
(Figure 3c and d).

Further evidence of the localized level of degradation of individual grains is provided by
comparing Figure 4b and 4d, illustrating two consecutive secondary electron hyperspectral
images of the same cross-section area of the biased sample in the energy range 0-6 eV
(additional images provided in SI Figure S3). In Figure 4d, 20-30 nm size dots were found to
appear at the surface of bright degraded grains (red circles, magnified image in Figure S4), but
not at the surface of darker grains. This is in agreement with other works showing that even
moderate e-beam currents (86 pA) and acceleration voltages (10 kV) can induce an aggressive
and very localized degradation of perovskite films due to ion migration. In comparison, the
reference sample did not show the appearance of e-beam induced-nano scale dots (Figure 4c).
These results confirm that the composition of some specific crystalline grains was modified by
biasing the device in humid air, and thus became even more sensitive to further degradation by the e-beam. We note that SEHI could be a fast and efficient novel method to investigate the chemical changes due to degradation in individual grains inside the film, which is particularly interesting for PSCs.

**Figure 5.** Integrated charge density (proportional to mobile ionic defects density) estimated from the integration of the transient part of the dark current at short-circuit condition just after biasing at 0.5 V, 0.85 V and 1 V for (a) a fresh perovskite device and (b) a device aged in air for two days.

To further understand the complementary roles of ion migration and humidity in the bias-induced degradation of perovskite devices, solid-state stepped chronoamperometry measurements were carried out. At this stage, the question is: can we correlate the degradation of
the device performance under electrical bias with the density of mobile ionic defects in the perovskite, and how is it related to the presence of moisture?

In Figure S5, dark current measurements for square wave voltage with various peak voltages show the “discharging” behavior below $V_{oc}$ and “charging” behavior above $V_{oc}$ until steady-state is reached after a few seconds. These transient behaviors have been previously explained by the movement of ions in the perovskite layer which screen the electric field for seconds as the dynamics response of ions is much slower than electrons and holes\textsuperscript{32,9}. If other mechanisms than ion migration may affect the current response, such as ferroelectric polarization or trapping of electrons at the interfaces, the rate of these mechanisms would be likely to be fast\textsuperscript{33,34} relative to the slow timescales of current response shown in Figure S5.

The transient response of the dark current to a potential step has been used to estimate the density of mobile ions present in perovskite devices. In Figure S6, square wave potential steps between 0 and 1 V were applied on the device while the dark current was monitored. The shape of the dark current curve can be explained similarly to transistor-type device used by D. Li \textit{et al.}\textsuperscript{35}: when a positive bias (1 V) is applied on the device, a current is immediately measured due to the drift of electronic carriers in the device. The bias then induces mobile ions to gradually move toward the contacts until equilibrium is reached after a few seconds when the ion-induced electric field cancels the applied bias in the perovskite bulk, which is observed by the stabilization of the dark current. After removing the applied bias (0 V), the previously accumulated ions at the perovskite–electrode interfaces create a net reversal electric field, leading to an instant switching of the polarity of the electronic current. The accumulated ions gradually diffuse away from the perovskite–electrode interfaces, resulting in the decrease of the ion-induced electric field and thus the reduction of the dark electronic current. If it is quite clear
that there is interplay between the ionic and electronic current, it may be asked the origin of
electronic carriers in the dark. It can be explained by the build-up of an electrostatic potential at
the interfaces due to the band bending induced by ions accumulation, as explained by R.
Gottesman et al.\textsuperscript{36}. When the bias is switched off, the electrostatic potential at the interface
remains for a few seconds before the device is “discharged” when mobile ions reached a new
equilibrium.

From this picture, we see that the magnitude of the ion-induced electric potential (or amount of
displaced mobile ions) will affect the magnitude and time constant of the measured dark current.
By integrating the area below the electronic current decay, we can thus estimate an “integrated
charge density” which is proportional to the amount of ions displaced in the device (mobile ions
density). This measurement is similar to a recent work by S. Bae et al.\textsuperscript{9} except that the dark
current is integrated after the voltage is removed.

Two cells from the same sample were measured shortly after fabrication and after storage in air
(40\% RH) for two days (Figure 5). For the fresh cell, the integrated charge density is $2 \times 10^{18}$
cm$^3$ at 0.5 V bias and increases up to $1.2 \times 10^{19}$ cm$^3$ at 1 V. For the aged cell, the charge density
is similar at 0.5 V bias ($3 \times 10^{18}$ cm$^3$) but considerably increases with higher voltage, up to $7.2 \times$
$10^{19}$ cm$^3$ at 1 V, which is 6 times higher than the fresh cell. As the sample was exposed to air
without encapsulation, environmental moisture entered within the perovskite film and likely
formed a hydrated phase containing more loosely bond ions.\textsuperscript{10} At 0.5 V forward bias, the electric
field is too low to activate these ions, which remain fixed in the crystalline structure. Thus,
similar charge density is measured for the fresh and aged cells. However, at 0.85 V or 1 V the
electric field is then sufficient to activate loosely bond ions and allow them to migrate in the
hydrated perovskite (aged cell), which explains the significant increase in charge density. These
results can be correlated with the degradation of the performance of the fresh and aged cells under dark bias. In Figure S7b, it is observed that the fresh cell degrades in hours instead of minutes for the aged cell, meaning that the perovskite devices having higher density of mobile ionic defects degrade faster under applied bias.

In addition, for the aged cell it can be observed that the charges density is slowly decreasing from $7.2 \times 10^{18}$ cm$^{-3}$ after the first potential step at 1 V to $5.4 \times 10^{18}$ cm$^{-3}$ after 4 potential steps. This trend can be explained as follow: as mobile ions move to the opposite interface under the effect of an electrical bias, they may be trapped if the space charge region at this interface is inverted, which happens if the applied voltage is greater than $V_{oc}$ or if a large amount of ions is displaced. As these ions are trapped, they cannot flow back in the perovskite bulk when the cell is switched back to short-circuit and do not participate in the measured current decay, which explains the charge density decrease at 1 V observed for the aged cell. This is not true for the fresh cell which has a higher $V_{oc} \geq 1$ V (current-voltage curves under 1 sun for the fresh and aged cells are shown in Figure S7a), i.e. similar to the applied bias, and lower charge density. In this case, the opposite interface is depleted but not inverted and mobile ions are less likely to be trapped. Indeed, the charge density remains unchanged even after four potential steps at 1 V.

Finally, the electrical bias stability of triple-cation mixed-halide perovskites was examined as they are known to have better moisture stability than conventional methylammonium lead halide perovskite$^{37,38}$. For this purpose, CS$_{0.1}$(MA$_{0.17}$FA$_{0.83}$)$_{0.9}$Pb(I$_{0.83}$Br$_{0.17}$)$_3$ perovskite solution was prepared following a recipe described in another work$^{39}$ and used as absorber layer in devices with the same architecture and interlayers as described above. The final device gave a maximum efficiency of 11.6%, which is closed to state-of-the-art performance for triple-cation perovskite processed on top of NiO$_x$.$^{40}$ Then, the dark bias stability of the device was measured by applying...
1 V forward bias for 1 h in 40% RH air. The J-V curves before and after aging are shown in Figure S9 and compared to a reference cell (without bias). The performances of the biased and non-biased cells are similar after aging, which indicates that the electrical bias has little or no effect on degradation for this kind of cells. This is ascribed to the better stability of Cs$_{0.1}$(MA$_{0.17}$FA$_{0.83}$)$_{0.9}$Pb(I$_{0.83}$Br$_{0.17}$)$_3$ perovskite to moisture. By avoiding the formation of a hydrated phase when exposed to moisture, the triple-cation mixed-halide perovskite device is unsensitive to the applied bias, which is on the contrary to CH$_3$NH$_3$PbI$_3$ cells which suffer from strong degradation when using the same aging conditions (Figure 1(a)). This is consistent with previous results showing that moisture can enhance ionic mobility in CH$_3$NH$_3$PbI$_3$, which will eventually have a detrimental and irreversible effect on the stability of the cells when an external bias is applied. Multiple-cation mixed-halide perovskites could be an efficient alternative for improving moisture and electrical bias stability of perovskite solar cells.

Conclusions

In summary, we showed that humidity is the main factor for the bias-induced degradation of inverted device based on CH$_3$NH$_3$PbI$_3$ perovskite in the dark, as degradation due to oxygen (superoxide) is not observed. Besides, XRD does not show any AgI, which implies that degradation related to Ag electrode is likely to be mild (if any), which is different to standard device structure, and can be explained by the fundamental difference that a positive bias is applied to the ITO in the inverted device structure (which attract I$^-$ to the ITO side). Our results suggest that, as the perovskite solar cell is exposed to a humid environment, moisture penetrates in the film preferentially from the edges of the cell and likely forms a hydrated perovskite phase which contains a high density of loosely bond ions. Under the effect of an external electric field, mobile ionic defects drift in the perovskite, which results in the accelerated degradation of the
active layer and eventually irreversible conversion of CH$_3$NH$_3$PbI$_3$ into PbI$_2$. Novel Raman –
transient photoelectrical technique is developed which allows probing local photoelectrical
properties on a highly degraded and non-degraded region. It indicates that the presence of a large
amount of PbI$_2$ suppresses or strongly reduces the migration of ionic species in the residual
perovskite. By applying advanced cross-sectional hyperspectral secondary electron imaging
technique, inhomogeneous degradation of distinct crystal grains was observed at the nano-scale
which could be explained by variations in the structural quality and ionic defects density from
grain to grain. Our study provides significant insight into the bias degradation of inverted
perovskite solar cells in the dark. Importantly, the new techniques developed and applied here
should be powerful to study the degradation of perovskite solar cells and other optoelectronic
devices in general.

Conflicts of interest
There are no conflicts to declare.

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Emitting Diodes Based on Solution-Processed Self-Organized Multiple Quantum Wells.  


The dark electrical bias degradation of inverted perovskite solar cells is due to ion migration in presence of moisture.