

# Review on Melt Flow Simulations for Thermoplastics and Their Fiber Reinforced Composites in Fused Deposition Modeling

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## Abstract

Fused deposition modeling (FDM) has been one of the most widely used additive manufacturing techniques due to its ease of use and widespread availability. Recent years have witnessed a renaissance of research interest in FDM due to its rapid advances in creating high-performance and industrial-grade parts, e.g., fiber-reinforced polymer composites. Considerable efforts have been made to model the FDM process to reveal the underlying mechanisms that are not yet well understood. This study aims to provide a comprehensive review of recent progress on the melt flow simulations of polymers and their fiber-reinforced composites in FDM. Specifically, analytical and numerical methods for modeling the polymer melt flow in the extrusion and deposition are summarized and discussed. Additionally, the FDM process simulation for short and continuous fiber-reinforced polymer composites as an emerging research field is outlined and discussed. Finally, the outlook of future work on the numerical simulation of FDM is provided.

**Keywords:** Process simulation; Fused deposition modeling; Melt flow; Fiber-reinforced composites; Computational fluid dynamics

## 1. Introduction

Additive manufacturing (AM), also called three-dimensional (3D) printing, is a method based on computer-aided design (CAD) models to build complex parts by material addition in a layer-wise manner. Compared with conventional manufacturing techniques, AM provides more degrees of freedom in design and manufacturing to create high-performance, complex

38 parts at low cost and high velocity. Due to continuous efforts, significant advances have  
39 been made to turn AM from rapid prototyping tools into a viable production option. Since the  
40 idea of AM was first introduced by Hull [1,2], various AM techniques have been developed,  
41 such as vat photopolymerization, material extrusion, binder jetting, material jetting, sheet  
42 lamination, powder bed fusion, and directed energy deposition [3].

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44 Among these techniques, fused deposition modeling (FDM), or fused filament fabrication  
45 (FFF), has been one of the most widely used AM techniques due to its ease of use and  
46 widespread availability. As a material extrusion technique, the FDM machine feeds  
47 continuous filaments into a heated liquefier and deposits these melt flows through pressure  
48 onto a print bed to build up 3D objects. Early FDM machines are mainly at the desktop level  
49 and used for fabricating engineering thermoplastics, e.g., polylactic acid (PLA),  
50 thermoplastic polyurethane (TPU), and polyamide (PA). The applications of parts fabricated  
51 by these machines are limited due to their poor mechanical properties and geometry  
52 accuracy. For a long period, the industrial value of FDM has been heavily underrated, and  
53 less research attention has been given to this technology in contrast to other AM techniques.  
54 However, the emergence of several innovative FDM techniques for creating high-  
55 performance parts has brought a renaissance of research interest to this field. For instance,  
56 specialized FDM techniques have been developed to process high-performance  
57 thermoplastic materials, e.g., Polyetheretherketone (PEEK) [4] and fiber-reinforced polymer  
58 composites (FRPC)[5]. The modulus and strength of FDM fabricated FRPC can be  
59 improved up to one or two orders of magnitude than using polymers alone. Also, big area  
60 additive manufacturing (BAAM), a variant of the FDM process, has been developed to  
61 increase the build rate and extend the print volume for large-scale parts [6,7]. The strong  
62 demand for these emerging techniques has motivated researchers to gain deep insight into  
63 the underlying mechanisms of the FDM process.

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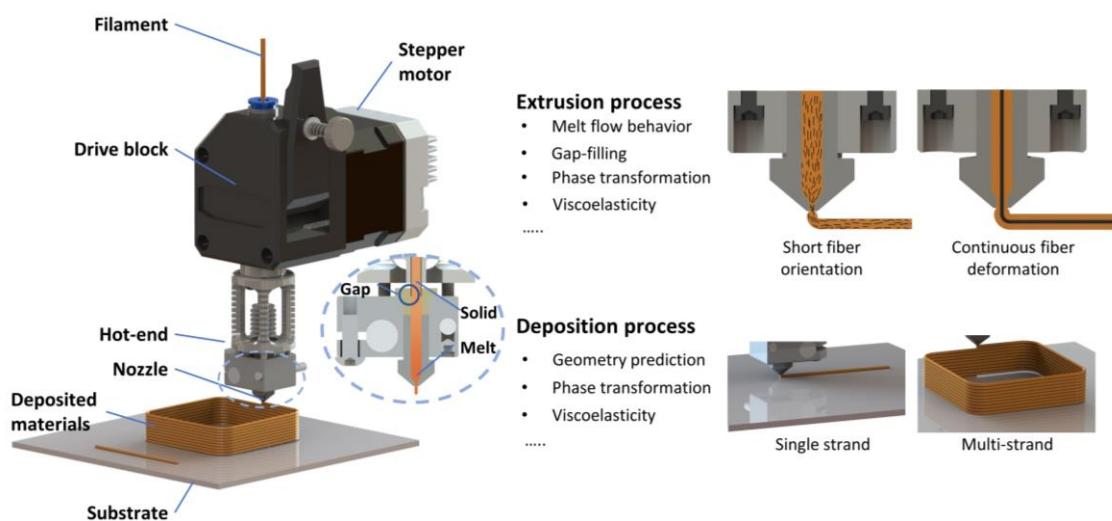
65 One critical mechanism is the polymer melt flow during the extrusion and deposition, which  
66 underpins the FDM technique and its advanced variants. The behavior of melt flow directly  
67 affects the FDM process's producibility, repeatability, and reproducibility. A clear  
68 understanding of this mechanism will form the basis for further applications, e.g., nozzle  
69 design and in-process control. However, most existing research [8–18] investigates the melt  
70 flow indirectly, in which its influences on the performance of fabricated parts are studied  
71 through experiments. In other words, parts are fabricated with different melt flows by  
72 adjusting the process settings, e.g., printing velocity [8,9,17] and nozzle temperature [9,13],  
73 thereafter their properties are characterized. However, the above methods cannot provide  
74 sufficient information to guide researchers for more complicated scenarios, e.g., extrusion  
75 at varied velocities and extrusion for extreme conditions. In-process monitoring is an  
76 attractive approach for directly observing the ongoing phenomenon [19–22]. However, the

77 limited accessibility makes the measurements prohibitively difficult, which is intangible to  
78 implement. Simulation-based methods provide a versatile solution to intuitively and rapidly  
79 understand the evolution of the melting and deposition process.

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81 This review aims to summarize the recent progress on the melt flow simulations of FDM for  
82 polymers and their fiber-reinforced composites. Fig.1 illustrates the fundamentals of the  
83 FDM process. The printhead, which is the core component of an FDM printer, consists of a  
84 filament feeding mechanism and a hot-end. The feeding mechanism, made up of a stepper  
85 motor and a drive block, continuously supplies the raw material filament into the hot-end.  
86 The hot end is responsible for heating the filament to be melted and successfully extruded,  
87 then, the extruded molten polymer brings into contact with the substrate and follows the  
88 printhead to move layer by layer to form the part. The influence of the deformation of the  
89 filament on the printing process can be ignored in the feeding process. Therefore, the  
90 simulations of the FDM process mainly focus on the process from the filament entry into the  
91 hot end to deposition on the substrate. The melt flow simulations can be divided into two  
92 parts: the extrusion process through a nozzle and the deposition process above a print bed.  
93 The thermodynamic models of pure polymer in the nozzle of FDM with analytical and  
94 numerical methods are discussed first. Then, an overview of numerical research about the  
95 polymer melt flow in the deposition process of FDM is presented examining both single-  
96 strand and multi-strand deposition. In addition, the FDM process simulation of short and  
97 continuous fiber-reinforced composites is also explored due to its significantly different  
98 phenomena. Finally, the article is concluded with an insight into the outlook of the future  
99 direction in numerical simulation of the FDM process. This work will contribute to exploring  
100 the application prospect of FDM process simulation and guiding the future research direction.

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Fig.1. Melt flow simulations for polymers and their fiber reinforced composites in fused deposition modeling.

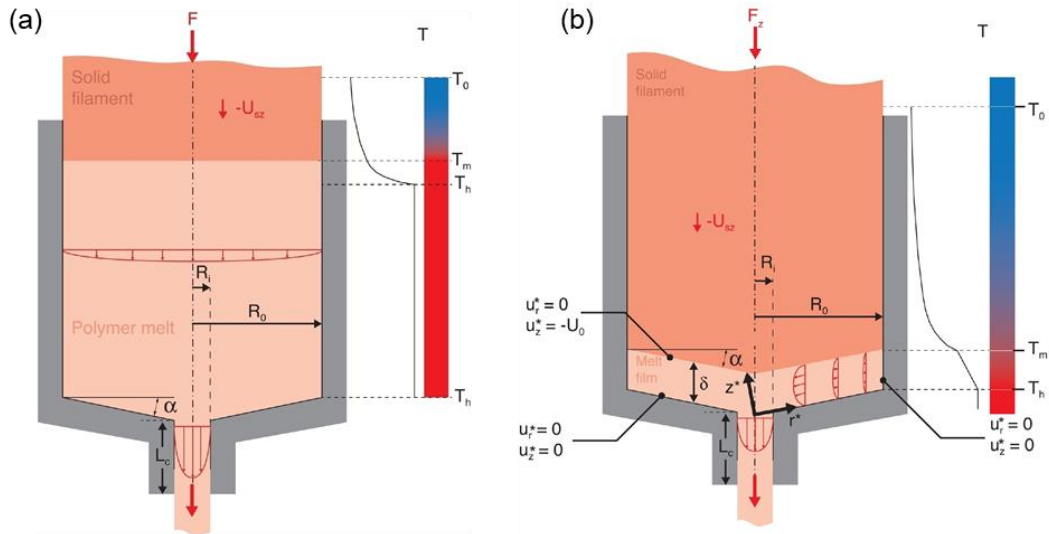
## 106 2. Intra-nozzle Process Simulations

### 107 2.1. Analytical models

108 Early attempts at intra-nozzle process simulations are mainly based on analytical melting  
109 models. Bellini et al. [23] are the first to utilize mathematic models as tools for investigating  
110 the melt flow phenomena and dynamics inside the liquefier. This research adopts a  
111 simplified non-isothermal flow in which the polymer in the hot-end is completely melted, and  
112 the temperature distribution is uniform and space-independent (Fig.2 (a)). The rheological  
113 behaviors of melted polymers, i.e., shear-thinning and temperature dependency, are  
114 described by the Power-law and Arrhenius models, respectively. The governing equations  
115 are based on the conservation law of fluid momentum and mass for calculating the pressure  
116 drop and evolution of flow rate. Meanwhile, the transfer function of the liquefier is  
117 established on the analogy of an amplifying circuit to predict the response of the flow field  
118 under different inputs, and its effectiveness is experimentally validated. The model is later  
119 applied as a part of a simulation system to support process optimization and real-time  
120 monitoring for feeding forces [24]. Phan et al. [25] further extend the model by considering  
121 the extensional viscosity effects using a modified Cogswell model [26] to investigate the  
122 pressure drop. The model's predictions on pressure drop have been compared with  
123 experimental measurements by adjusting the Trouton ratio, which is the ratio of extensional  
124 viscosity to shear viscosity and measures the strength of elasticity in the melt flow.

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126 The above models are not valid for processes with high feed rates, in which the filament is  
127 not completely melted until the bottom of the liquefier. Thus, Osswald et al. [27] have  
128 proposed an analytical model, in which a melt film is modeled on the surface of the conical  
129 section, for incorporating the phase transformation of polymer from solid to melt states (Fig.2  
130 (b)), and validated it with an in-house developed test setup. In later studies, Peng et al. [28]  
131 have established an analytical model to describe the flow profile inside the hot-end channel  
132 with assumptions of isothermal flow and Power-law fluid and applied Computer Tomography  
133 (CT) to scan the nozzle with a solidified pigment filament. It is found that the model prediction  
134 deviates from the measured velocity profile due to the oversimplification of the melt flow's  
135 non-uniform temperature distribution. In addition, a more advanced analytical model has  
136 been proposed for studying the deformation of the polymer during the FDM process [29].  
137 They have adopted a modified continuum molecularly aware polymer model with flow-  
138 induced disentanglement and found that the polymer deformation is affected primarily by  
139 the deposition flow rather than the intra-nozzle flow. Also, the polymer deformation  
140 increases with the printing velocity.



141

142 Fig.2. (a) Melting and flow model assumed [23] by Bellini et al. [23]; where  $T_0$ ,  $T_m$ ,  $T_h$  are initial  
 143 temperature, melt temperature of polymer, and heating temperature from hot-end. The  
 144 temperature distribution of polymer in the hot-end is uniform; (b) Improved model in which a  
 145 small melt film on the surface of conical section plays a significant role in phase transformation.

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## 156 2.2. Numerical models

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Although previous studies have investigated various numerical melting models, the following standard procedure can be summarized to construct computational fluid dynamics (CFD) models. The modeling process includes nozzle geometry modeling, fluid flow model selection, material property definition, boundary condition definition, and numerical methods selection. The selection of fluid flow models, i.e., laminar or turbulent flow models, depends on the Reynolds number. In a typical scenario, the outlet diameter of the nozzle, outlet velocity, and polymer density, are less than 1 mm, 120 mm/s, and 10 g/cm<sup>3</sup>, respectively. Thus, the Reynolds number is less than one and the fluid flow is laminar. In terms of material property definition, the basic thermophysical and rheological properties of materials that should be considered are density, viscosity, specific heat, thermal conductivity, melting enthalpy, and melting temperature. Afterward, boundary conditions, including the velocity at

168 the inlet, atmospheric pressure at the outlet, wall temperature, and heat transfer conditions,  
169 need to be defined. Once these conditions are established, numerical methods such as the  
170 finite difference method, finite element method (FEM), and finite volume method (FVM) are  
171 used for solving the governing equations. These simulations are often performed within CFD  
172 software tools, including Fluent, Polyflow, OpenFOAM, COMSOL, and FLOW-3D [4,30–39].  
173  
174 Previous simulations have investigated the use of various constitutive models for  
175 thermoplastic melts. Table 1 summarizes typical thermophysical and rheological models. In  
176 terms of material density, temperature-dependent models have been studied [4,31,33,38],  
177 though most research treats the density as constant. Also, the impact of material  
178 compressibility on the density has been only considered in the study of Kattinger et al. [38].  
179 Regarding the rheology behavior of polymer melts, most studies have utilized classical  
180 Generalized Newtonian Fluid (GNF) models [4,31–35,37–40], such as the Carreau-Yasuda,  
181 Power-Law, and Cross models, to describe the shear rate ( $\dot{\gamma}$ ) dependent behavior of the  
182 viscosity. Meanwhile, Arrhenius [32,35,37,39,40] or Williams-Landel-Ferry (WLF)  
183 [4,33,38,41] models are often adopted for depicting the temperature dependence of  
184 viscosity. Meanwhile, some research has applied the viscoelastic model, e.g., Giesekus  
185 [30,41] and Phan-Thien-Tanner (PTT) [34] models, which consider the elastic effect and are  
186 more accurate yet complex, compared to traditional rheological models. For thermal  
187 properties, previous studies [31,32,34,38] have modeled the specific heat capacity as a  
188 function of temperature. Moreover, recent studies [33,38] have taken into account other  
189 complex material properties, e.g., the pressure-dependent glass transition  
190 temperature  $T_{glass}(p)$ , temperature-dependent melting enthalpy  $h(T)$ , and temperature-  
191 dependent thermal conductivity  $\kappa(T)$ , and temperature-dependent thermal radiation  
192 intensity  $\dot{q}_{rad}(T)$ .  
193

Table 1 Summary of numerical melting models for the FDM process.

Author	Research application	Material	Parameters of the thermal flow model			Tools	Ref.
			Density	Viscosity	Thermal		
Shadvar et.al	Melt flow behavior	Acrylonitrile Butadiene Styrene (ABS)	Constant	$\mu(\dot{\gamma}, T)$ (Power-law& Arrhenius)	Constant	Fluent	[37]
Phan et.al	Melt flow behavior & viscoelastic model	Poly lactide (PLA)	Constant	$\mu(\dot{\gamma})$ (Carreau-Yasuda) $/\mu(\dot{\gamma}, \sigma)$ (PTT)	$C_p(T)$ , others are constant	Fluent/ Polyflow	[34]
Ufodike et.al	Melt flow behavior	ABS	Constant but differ in phases	$\mu(\dot{\gamma}, T)$ (Power-law& Arrhenius)	$C_p(T)$ , others are constant	Fluent	[32]
Kattinger et.al	Melt flow behavior	Polystyrene (PS)	$\rho(p, T)$	$\mu(\dot{\gamma}, T)$ (Cross&WLF)	$C_p(T), \kappa(T)$ $h(T), T_{glass}(p)$	OpenFOAM	[38]
Schuller et.al	Melt flow behavior	Polycarbonate (PC)	Constant	$\mu(\dot{\gamma}, \sigma)$ (Giesekus)	-	OpenFOAM	[30]
Serdeczny et.al	Melt flow behavior & geometric effect	ABS	Constant	$\mu(\dot{\gamma}, \sigma, T)$ (Giesekus& WLF)	Constant	OpenFOAM	[41]
Serdeczny et.al	Melt flow behavior & gap-filling	PLA	$\rho(T)$	$\mu(\dot{\gamma}, T)$ (Power-law& WLF)	$\dot{q}_{rad}(T)$ , others are constant	FLOW-3D	[33]
Marion et.al	Melt flow behavior & gap-filling	ABS	Constant	$\mu(\dot{\gamma}, T)$ (Carreau-Yasuda & Arrhenius)	Constant	Self-programming	[40]
Zhang et.al	Temperature distribution & change	PLA	$\rho(T)$	$\mu(\dot{\gamma})$ (Power-law)	$C_p(T)$ , others are constant	Fluent	[31]
Idris et. al	Temperature setting	PLA	Constant	$\mu(\dot{\gamma}, T)$ (Power-law& Arrhenius)	Constant	Fluent	[39]
Papon et.al	Nozzle outlet design	PLA	Constant	Constant	Constant	Fluent	[36]
Go et.al	Nozzle design & velocity bound setting	-	Constant	$\mu(\dot{\gamma}, T)$ (Power-law& Arrhenius)	Constant	COMSOL	[35]
Wang et.al	Temperature & velocity bound setting	Polyetheretherketone (PEEK)	Constant	$\mu(\dot{\gamma}, T)$ (Cross-WLF)	Constant	Fluent	[4]

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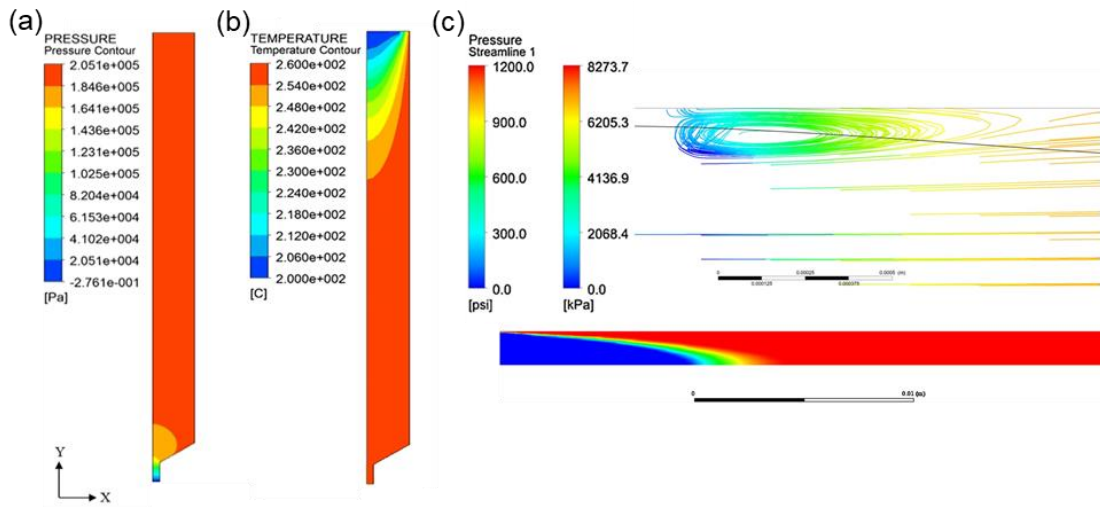
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Although the material models chosen in previous studies are different, some similarities can be found among the results obtained [4,30–39]. Firstly, most studies have found that the pressure drop changes mainly occurred in the conical part of the nozzle, and the magnitude of the pressure drop is proportional to the polymer feed rate (Fig.3 (a)). Secondly, the temperature distribution of the molten polymer, and consequently the solid-melt interface, in the hot-end channel is a smooth cone-like surface with the tip pointing downward (Fig.3 (b)). The position of the interface moves toward the nozzle outlet with an increase in the feed rate. Thus, the polymer filament is prone to clog the nozzle at a high feed rate due to insufficient melting. In addition, the recirculation vortex is found between the liquefier wall and the filament that just gets molten, as depicted in Fig.3. (b). Some studies have speculated that this recirculation vortex is formed by the drag flow generated by the moving

207 filament that acts to counteract overflow [32–34]. In terms of heat transfer conditions, some  
208 studies [33,34,42] have found that the increase in the wall heat transfer coefficient and  
209 heating temperature accelerates the melting of the polymer.



210

211 Fig.3. Pressure (a) and temperature (b) distribution of the ABS extrusion. Reprinted with  
212 permission from reference [37], Copyright 2019, Springer. (c) Above: the distribution of the  
213 pressure drop in the vortex region; below: the region of solidification (blue) and melting (red).

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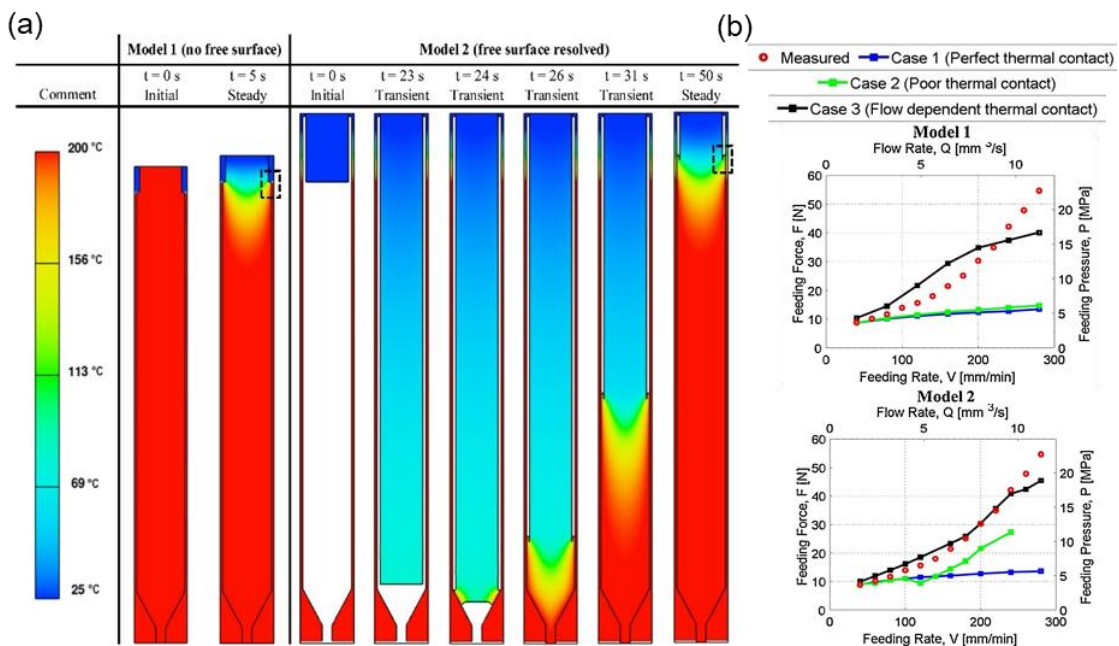
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216 The aforementioned simulations consider the polymer melt as GNF, which neglects the  
217 elastic memory effect of the polymer. Thus, applying the viscoelastic model to examine the  
218 elastic effect on polymer during the FDM process is of great interest. Tomás Schuller et al.  
219 [30] have simplified the fluid domain from 3D to 2D (two-dimensional) according to the axial  
220 symmetry of the FDM nozzle pipeline, which significantly reduces computational time and  
221 cost. The research has assumed the flow in the nozzle as isothermal flow and adopted the  
222 Giesekus model to characterize the polymer viscoelasticity. It is revealed that whether the  
223 tensile flow is stronger than the shear flow at a certain point in the fluid domain depends on  
224 the positive or negative normal stress difference at that point. There is a large interaction  
225 force between the polymer melt and the nozzle's conical part, which changes the  
226 unidirectional stress state of the melt and its flow direction, enlarge the normal stress  
227 difference, and strengthen the elastic effect. The shear elastic stress is dominant in the  
228 backflow in the solid filament that enters the print-core and the liquefier wall. The study has  
229 also predicted the location of the recirculation vortex that occurred. Serdeczny et al. [41]  
230 have improved this model to reflect the actual non-isothermal flow in the FDM process by  
231 employing the WLF model for depicting the temperature dependency of the rheological  
232 parameters in the Giesekus model. In contrast to the commonly used GNF model, the use  
233 of this viscoelastic model can significantly improve the prediction accuracy of the pressure  
234 drop for printing with high feed rates.



235

236 In recent studies, the filling process for the gap between the polymer filament and the inner  
237 wall of the hot-end channel has been investigated. Serdeczny et al. [33] and Marion et al.  
238 [40] have improved the geometric model and applied a multiphase model to characterize  
239 the gap-filling process. Fig.4 (a) shows the comparison between models without and with  
240 the gap-filling process, which indicates more time is required to reach the steady state to  
241 form a stable backflow region due to the gap-filling process. However, it is found that when  
242 the feed rate is higher than a certain limit, the position of the backflow region in the new  
243 model will not remain stable. Otherwise, there is no difference between the two models.  
244 Researchers have speculated that this might be related to overflowing [33]. At high feed  
245 rates, the radial heat conduction will not be enough to melt the polymer in time, and the  
246 filament may be directly pushed down in the channel. When the filament contacts the conical  
247 section of the nozzle, the pressure applied to the polymer becomes larger. Then, the  
248 backflow is generated to fill the gap upward and increase the contact area between the  
249 filament and liquefier wall, which increases the melt temperature and decreases the  
250 viscosity and overall pressure. Such fluctuation of the gap-filling level repeats while the  
251 filament is feeding. Meanwhile, Ye et al. [22] have observed these phenomena using a  
252 transparent nozzle. In addition, it is found that the gap-filling level fluctuation occurs when  
253 the feed rate is beyond a threshold. A too-high feed rate will cause overflow from the hot-  
254 end and clog the print head. Meanwhile, the change of the feeding force under high feed  
255 rates has been studied by Serdeczny et al. [33], as depicted in Fig.4. (b), which has revealed  
256 that considering the gap-filling process and a proper thermal contact condition can improve  
257 the simulation accuracy.



258

259 Fig.4 (a) Comparison of models without (Model 1) and with (Model 2) gap-filling process; and  
260 (b) the change of the feeding force under different models and thermal contact conditions.

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263 **2.3. Applications of numerical method in FDM**

264 Through the above numerical simulations, the influence of the hot-end's material, geometry,  
 265 and heat transfer conditions on printing performance, such as maximum printing velocity  
 266 and heating temperature, can be predicted. Thus, there are some potential applications in  
 267 engineering design for FDM numerical simulation, such as optimizing the parameters for  
 268 the printing process, planning the printing path, or assisting with the print head design  
 269 [31,35,36,39,43]. For example, Idris et al. [39] have studied the solid-liquid distribution of  
 270 the polymer in the liquefier to determine the heating temperature of the printer hot-end.

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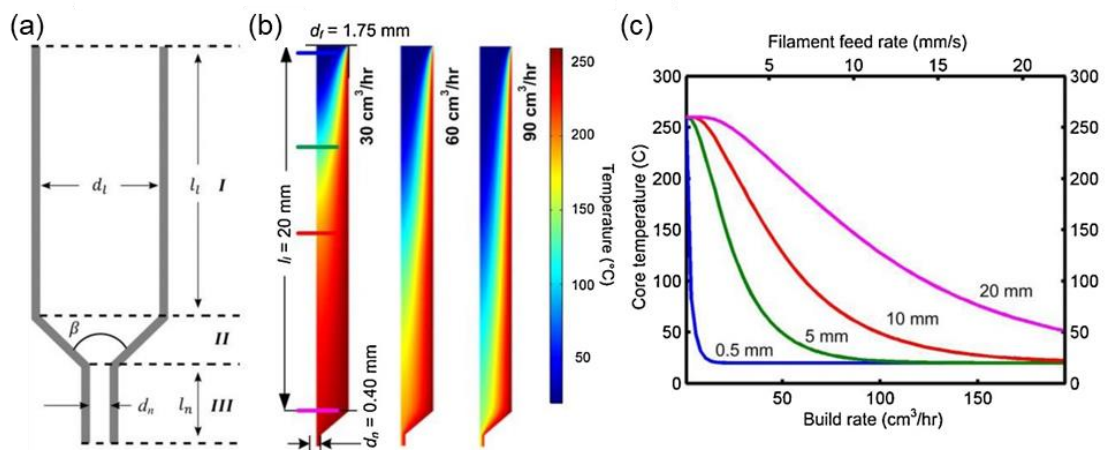
272 As shown in Fig. 5, some studies have determined the range of the feed rate according to  
 273 whether the polymer is melted sufficiently before the nozzle outlet, which ensures fluent  
 274 extrusion for the polymer [31,35]. Moreover, Comminal et al. [43] have concluded the  
 275 characters of deposition at the corner and applied it in motion planning by numerical  
 276 simulations.

277

278 In terms of print head design, numerical simulation is generally applied for verifying or  
 279 comparing its printing performance. For instance, Papon et al. [36] have investigated the  
 280 influence of different shapes of the nozzle outlet on printing performance by comparing  
 281 pressure drop at the same feed rate. Go et al. [34] have used numerical simulations to verify  
 282 the designed liquefier by simulating whether the polymer could be extruded smoothly. In  
 283 addition, the numerical model has been persistently improved for investigating the change  
 284 of feeding force with feed rate increasing [33,38,41], which is one of the important bases for  
 285 print head design under various printing requirements, e.g., FDM rapid printing, conformal  
 286 printing, printing for high-performance materials (e.g., PEEK) or soft materials (e.g., TPU).

287

288



289 Fig. 5. A representative liquefier design case: First step: (a) design dimensions of liquefier;  
290 Second step: numerical simulation to determine the process parameters, (b) simulate the  
291 temperature distribution within the hot-end channel under different feed rate to ensure polymer  
292 melt sufficiently and easily extruded, (c) polymer temperature versus feed rate at noted  
293 distances in (b) to determine the range of feed rate. Reprinted with permission from reference  
294 [35], Copyright 2017, Elsevier.  
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### 296 **3. Deposition Process Simulations**

297 In the FDM process, the polymer melt is extruded through the hot-end and then deposited  
298 onto the platform to form parts layer by layer. The phase transformation of the polymer melt  
299 flow within the deposition process is a decisive factor that affects the deposition quality.  
300 Modeling such a process is a nontrivial task since it includes the forming of strands when  
301 the melt leaves from the nozzle and the partial re-melting to re-solidifying of existing strands  
302 near the strands being extruded. In addition, the deposition process simulation is a multi-  
303 phase problem with high computational costs, which involves the time-dependent phase  
304 change and the intricate interplay between different phases. The multi-phase model should  
305 be able to capture the evolution of the interface between air and polymer melt, i.e., the  
306 forming process of strands. Furthermore, previous studies [44–47] have found the shape of  
307 the extruded strand directly affects the performance of printed parts, e.g., surface roughness  
308 and porosity. Hence, recent studies have extensively investigated the use of multi-phase  
309 models for predicting the geometry of the strands. The rest of this section discusses the  
310 modeling of the two consecutive processes during the deposition. One is the single strand  
311 deposition, i.e., the strand-forming process when the melt leaves from the nozzle; another  
312 is multi-strand deposition corresponding to the re-solidification of existing strands near the  
313 strand being extruded.

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#### 315 **3.1. Single strand deposition**

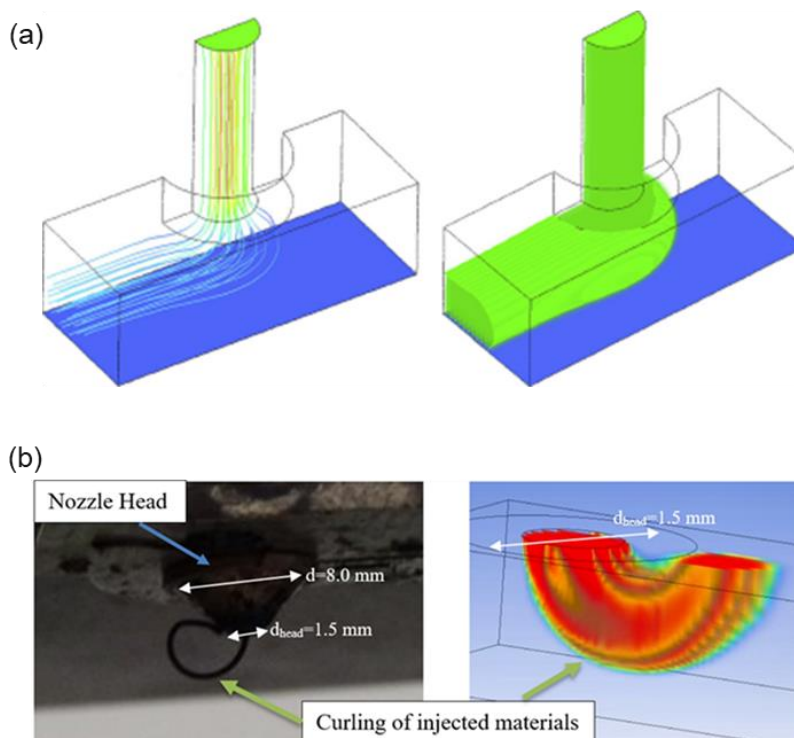
316 Single-strand deposition means that polymer melt extrudes from the FDM nozzle to form a  
317 single bead on the print bed, as illustrated in Fig.6 (a). Numerical models with various  
318 complexities and accuracies have been developed, including models based on Newtonian  
319 [48] or non-Newtonian fluids [49,50], including isothermal [48,51–54] or non-isothermal  
320 flows [41,42]. For non-Newtonian fluid assumption, the Cross-WLF or Cross model [49,50],  
321 has been adopted. Meanwhile, material thermal parameters and heat transfer conditions  
322 [49,50], e.g., specific heat capacity, phase transition temperature, and wall heat transfer  
323 coefficient, have been considered for non-isothermal flow. It is found that the shear rate,  
324 temperature-dependent viscosity, and heat transfer conditions have a significant impact on  
325 improving simulation accuracy [50]. Meanwhile, the effects of surface tension and buoyancy

326 on extrusion deposition have also been considered in previous works [49,52].

327

328 In the studies above, the influences of key factors, e.g., normalized printing head velocity  
329  $V/U$  (extrusion volumetric flux  $U$  and printing head velocity  $V$ ) and the normalized gap  
330 height  $g/D$  (nozzle position  $g$  above the substrate and the nozzle diameter  $D$ ), on the cross-  
331 sectional profile of the strand have been studied [48,49,52–54]. The results show that the  
332 thickness of the strand increases with the increase of  $g/D$ , while curling phenomena [46]  
333 (Fig.6 (b)) can be observed if the normalized gap height  $g/D$  exceeds a threshold value. In  
334 addition, at the same  $g/D$ , with the increase of  $V/U$ , the ratio of thickness to the width of  
335 the strand and printing force will decrease. The simulation results match well with the  
336 experimental measurement at a low  $V/U$  and  $g/D$  [53]. Meanwhile, regarding the levelness  
337 of the strand, the simulation results deviate from the experimental measurements under the  
338 fluctuating feeding force. This is caused by the neglect of the viscoelastic effect within the  
339 model [52]. Furthermore, previous studies have also simulated the deposition process of  
340 core-shell polymer manufacturing [54], which is a variant of FDM that processes two  
341 polymers with distinct glass transition temperatures. It has been found that a good packing  
342 effect and a high-volume fraction of the core can be achieved by using large  $V/U$ ,  $g/D$ , and  
343 inner nozzle diameter.

344



345

346 Fig.6. (a) An example of simulation for the deposition process of the single strand. Reprinted  
347 with permission from reference [48], Copyright 2018, Elsevier. (b) Curling phenomena

348 observed in the experiments and numerical models. Reprinted with permission from reference

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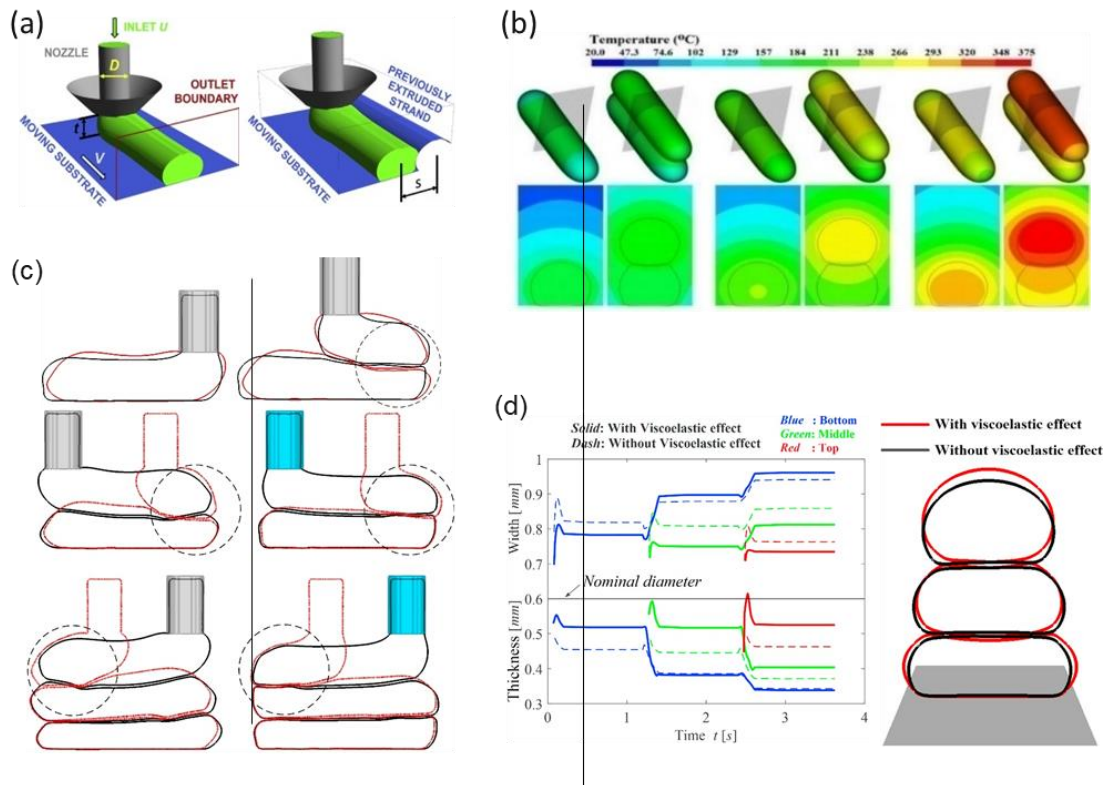
### 351 **3.2. Multi-strand Deposition**

352 In multi-strand deposition simulations, previous studies have focused on predicting the  
353 interface between strands, which is critical to the geometry and mechanical properties of  
354 the final part. There have been a few existing methods for modeling and simulating the multi-  
355 strand deposition. The applied methods can be divided into three types according to the  
356 thermal condition and material constitutive models.

357

358 In Method 1, an isothermal flow assumption is made, the previously deposited and solidified  
359 material are treated as rigid bodies, and the last strand is deposited at the new substrate  
360 combined with previous solid strands and substrate [55] (Fig. 7 (a)). Method 2 only considers  
361 the non-isothermal assumption and assumes the interface between strands will change  
362 through a partial re-melting and re-solidification process [56,57] (Fig. 7 (b)). Based on  
363 Method 2, Method 3 employs the viscoelastic model [58,59] in numerical analysis and  
364 investigates the influence of viscoelasticity on the deformation between the strands [60]. It  
365 has been found that compared with the non-Newtonian viscous fluid model, there are  
366 obvious differences in extruding diameter, layer thickness, strand profile, and shape  
367 recovery over time (Fig. 7(c-d)). In other words, with the improvement of the model by these  
368 three methods, the prediction accuracy of multi-layer or multi-strand deposition will be  
369 gradually improved.

370



371

372 Fig. 7. (a) Interface between strands only determined by the previous strand surface. Reprinted

373 with permission from reference [55], Copyright 2019, Elsevier. (b) The interface between

374 strands changes with the reheating of the previous strand. Reprinted with permission from

375 reference [56], Copyright 2017, Emerald. (c) The interface changes during the deposition

376 process from the simulations with viscoelastic effect (black) and without viscoelastic effect

377 (red); and (d) thickness and width of each layer versus time from simulations with viscoelastic

378 effect (solid line) and without viscoelastic effect (dash line). Reprinted with permission from

379 reference [60], Copyright 2018, Elsevier.

380

#### 381 4. Process Simulations for Fiber-Reinforced Polymer Composites

382 FDM for FRPC is an emerging technique to fabricate lightweight but strong polymer

383 composites in which fibers are utilized as the reinforced phase. Typical fiber types include

384 carbon, glass, and Kevlar fibers. These fibers are either in short or continuous forms

385 corresponding to two different processes. For short fibers, their polymer composite filaments

386 can be directly used on commercial FDM machines without significant changes. However,

387 the weight contents of short fiber in composites are often limited due to the reduced fracture

388 strain. For continuous fibers, a co-extrusion strategy is often adopted to impregnate the fiber

389 bundles with molten polymers [61]. These fibers form continuous load paths within

390 fabricated composites, which possess one or two magnitudes higher mechanical properties

391 than short-fiber reinforced composites. Since fibers are highly anisotropic materials, the

392 alignment and orientation of fibers are critical to the mechanical properties of fabricated



393 parts. In addition, other characteristics, such as the interface bonding properties, play an  
394 important role in it. Therefore, investigating the material flow and its interactions with fibers  
395 during the extrusion and deposition processes through simulations will be helpful for refining  
396 the FDM process for FRPC.

397

#### 398 **4.1. Process simulations for short fiber-reinforced composites**

399 The melt flow simulation of short fiber composites is a fluid-structure coupling problem. In  
400 this problem, the fiber and molten polymer are considered as the solid and fluid phases,  
401 respectively. These two phases interact and influence each other. Accurate modeling of  
402 these complex interactions is one of the main challenges for process simulation. Table 2  
403 summarizes details of previous studies on process simulations for short fiber-reinforced  
404 composites. The simulations based on CFD methods, such as FVM [62,63], FEM, and  
405 particle projection method, are often used for simulating the flow field. Moreover, the  
406 approaches for modeling fibers could be divided into three types. Method 1 models the fibers  
407 as rigid rods with an aspect ratio far greater than 1 [64–75]. Method 2 discretizes the short  
408 fibers into several rigidly connected small spheres or particles in the flow field [62,76].  
409 Method 3 is to equivalent the mix of the fiber and polymer melt as pure fluid [77] and applies  
410 the moving particle semi-implicit (MPS) method to discretize the fluid into particles, in which  
411 the effect of fibers in the equivalent fluid is reflected in the viscosity model. In addition, this  
412 method has been validated by the experiment in predicting the cross-sectional profile of the  
413 deposited strand [77].

414

415 Table 2 Summary of numerical simulations for short fiber-reinforced composites in the FDM  
 416 process.

Author	Research application	CFD method	Fibers model	Degree of coupling	Isothermal or Non-isothermal	Ref.
Yang et.al	Fracture	FVM	-	-	Isothermal	[63]
Imaeda et.al	Deposition	MPS	-	-	Non-isothermal	[77]
Smith et.al	Fiber orientation	FEM	Orientation tensor	Weak	Isothermal	[66–70]
				Full		[64,65,71,75]
Bertevas et.al	Fiber orientation	SPH	Orientation tensor	Full	Isothermal	[72]
					Non-isothermal	[73,74]
Yang et.al	Fiber orientation & nozzle clogging	SPH	DEM	Full	Isothermal	[76]
		FVM				[62]

417  
418

419 Method 1 considers the fibers as rigid rods based on the Folgar-Tucker orientation tensor  
 420 evaluation model [78], in which the aspect ratio and volume fraction of fibers are considered.  
 421 With this method, the orientation of fibers in melt flow under the prescribed fluid velocity can  
 422 be simulated [64–75]. In addition, the relationship between fibers and melt flow can be  
 423 modeled as (i) weak coupling [66–70] or (ii) full coupling [64,65,71–75]. Specifically, the  
 424 weak coupling model neglects the presence of fibers in the computation of melt flow  
 425 kinematics. In the simulation, conservation equations of pure polymer melt flow are solved,  
 426 then the velocity distribution is substituted into the orientation tensor evaluation model to  
 427 get the fiber orientation state. This model is developed for studying the orientation of short  
 428 fibers [66]. It is found that the fiber orientation tends to align along the direction of polymer  
 429 melt flow, but the nozzle die swells hinder this trend. In addition, studies on BAAM have  
 430 revealed that the planar extrudate swell, different polymer constitutive models, and screw  
 431 swirling also have a direct impact on fiber orientation during the deposition process [67–70].  
 432 In contrast, the full coupling model considers the influence of fibers on the polymer melt flow,  
 433 where the fiber orientation tensor is coupled into the momentum equation and energy  
 434 equation in the form of stress (if a non-isothermal flow is assumed). Smith et al. [64] have  
 435 developed a full coupling model with FEM [65,71] and found that the average value of fiber  
 436 length distribution has a significant effect on the mechanical properties of the extruded  
 437 strand [62]. Bertevas et al. [72] have established a full coupling model through the  
 438 smoothed-particle hydrodynamics (SPH) method, a particle projection method, which  
 439 discretizes the fluid domain into a series of particles for solving complex fluid problems [70].  
 440 Due to its computational rapidity distinguished from other numerical methods (such as FVM)  
 441 for the CFD model, it is affordable to apply the non-isothermal conditions in the later  
 442 simulation [73,74]. Research has found that the degree of fiber alignment on the surface of  
 443 the deposited strand is better than that in the interior. This effect will become more

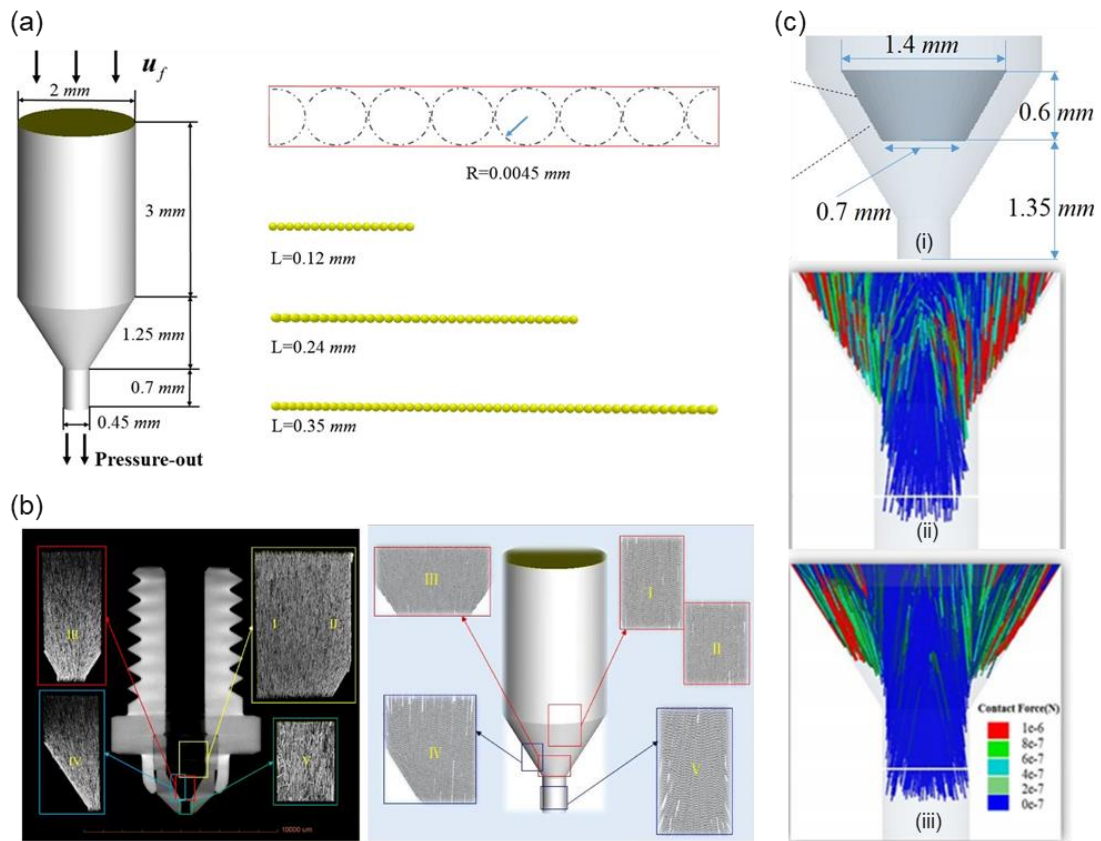


444 pronounced with the increase of fiber concentration and the ratio of feed rate  $u$  to printing  
445 head velocity  $V$  [72]. The effect of the thermal conductivity of fibers has been revealed to  
446 enhance the fiber alignment along the deposition direction on the top half of the strand but  
447 weaken the fiber alignment on the bottom half of the strand. Meanwhile, it has been also  
448 found that the extrusion and deposition of the current layer will randomize the fiber  
449 orientation within the previously deposited layer. This effect can be aggravated by the  
450 increase of fiber concentration and fiber aspect ratio [73,74].

451

452 Method 2, called the discrete phase model (DEM), discretizes the short fibers as several  
453 rigidly connected particles in the polymer flow (Fig.8 (a)). In contrast to other methods, DEM  
454 considers the deformation of the fibers. This method has been used to investigate the fiber  
455 alignment and clogging during the extrusion in the studies by Yang et al. [62,76], in which  
456 the SPH and FVM are applied to simulate the flow field. Moreover, the prediction of this  
457 method has been validated by CT images (Fig.8 (b)). Their studies have found that in the  
458 cone section of the nozzle, the fiber orientation tends to be parallel to the tapered wall, and  
459 the continued downward flow of the fibers causes cross-links. The longer the fibers are, the  
460 greater the possibility of cross-linking is. The cross-linking of the fibers results in a large  
461 pressure drop and even clogs the nozzle. One possible solution is to install a cone sleeve  
462 in the cone section of the nozzle, as illustrated in Fig.8 (c). Due to the cone sleeve, the  
463 reduction of the contact force between fibers near the outlet has been found in the simulation,  
464 which contributes to restrain the fiber cross-linking and nozzle clogging.

465

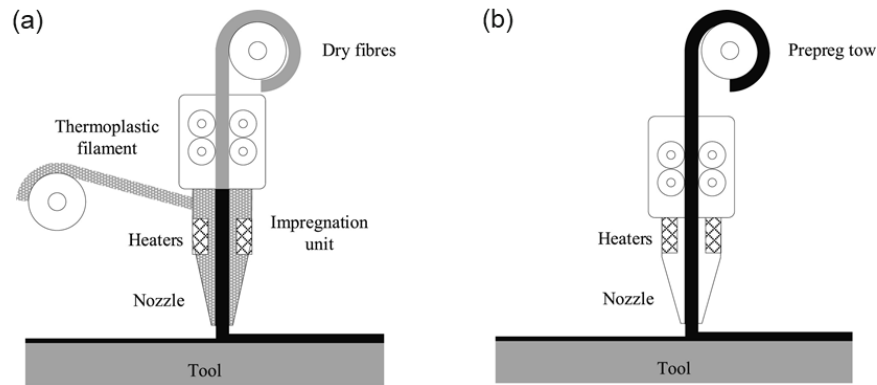


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Fig.8. (a) Modeling of fiber and nozzle, where fibers are constituted of discrete DEM particles. (b) Fiber distribution in five regions: CT images (left) and numerical results (right); (c) Sketch of the cone sleeve (i) and distribution of DEM particles colored by their contact forces for origin nozzle (ii) or using cone sleeve (iii). Reprinted with permission from reference [62], Copyright 2021, Springer.

#### 474 4.2. Process for continuous fiber-reinforced composite

475 As shown in Fig. 9, FDM for continuous fiber-reinforced composites (CFRP) can process  
 476 both dry and pre-impregnated fiber bundles consisting of hundreds to thousands of  
 477 continuous fibers. The critical difference is the existence of an in-situ impregnation process  
 478 during extrusion, which significantly influences the mechanical properties of fabricated  
 479 composites [79].

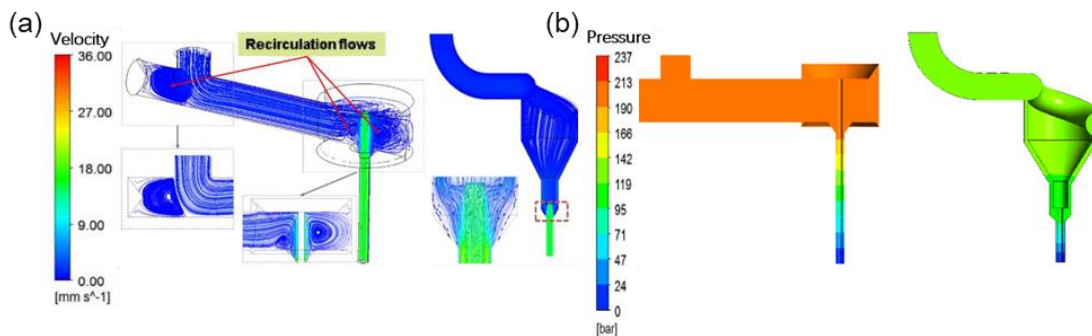


480

481 Fig. 9. Two typical FDM processes for CFRP: (a) in-situ impregnation and (b) pre-  
 482 impregnation. Reprinted with permission from reference [79], Copyright 2021, Elsevier.

483

484 There are limited studies on the simulating continuous fiber-reinforced thermoplastic melt.  
 485 In these studies, the fiber tow is often considered as a rigid sliding surface. Under this  
 486 assumption, Han et al. [80] have used a numerical method for solving the isothermal  
 487 Newtonian fluid model to calculate the pressure on the fiber to verify the feasibility of nozzle  
 488 designs. Albrecht et al. [81] have proposed a computational fluid dynamics model with the  
 489 non-isothermal viscosity model (Cross-Arrhenius) and viscoelastic model (K-BKZ), to assist  
 490 the design for the geometry of the hot-end. The velocity and pressure drop profiles shown  
 491 in Fig. 10 demonstrate the superiority of the modified hot-end design with fewer vortices,  
 492 residence times, and lower inlet pressure drops. In addition, the dimensions of the hot-end,  
 493 i.e., length and outlet diameter, has been optimized using design objectives including the  
 494 pressure drop, die swell ratio, and shear stress.



495

496 Fig. 10. (a) Streamlines within the 1st generation hot-end (left) and modified hot-end (right);  
 497 and (b) pressure distribution within 1st generation hot-end (left) and modified hot-end (right).

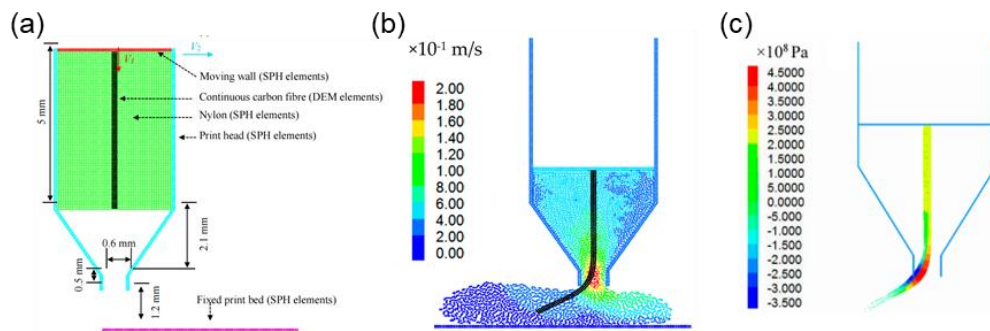
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Reprinted with permission from reference [81], Copyright 2019, Sim-AM 2019.

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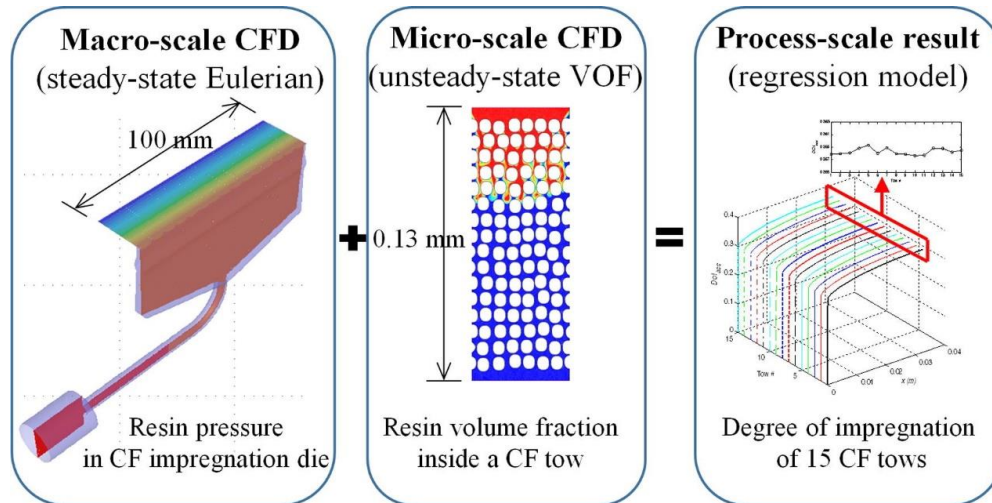
500 In addition, Yang et al. [76] have adopted the DEM to study the deformation of long fiber tow  
 501 within the melt flow by assuming the fiber tow as a line formed by a rigid connection of  
 502 particles (Fig. 11 (a)). Meanwhile, the polymer fluid is considered as the Newtonian fluid

503 under isothermal conditions. This study is the first to numerically investigate the extrusion  
 504 and the deposition process for CFRP-FDM (Fig. 11 (b)), considering the fiber deformation.  
 505 As shown in Fig. 11 (c), the shear stress concentration is conspicuous on the fiber at the  
 506 nozzle outlet due to the bending. As an improvement, this study has pointed out that the 2D  
 507 model can be extended to a 3D model in the future to study how continuous fiber is wrapped  
 508 by the resin during the deposition process.  
 509



510  
 511 Fig. 11. (a) Schematic diagram of numerical modeling for CFRP-FDM. (b) Simulation result of  
 512 CFRP-FDM extrusion and deposition process at  $t = 0.1$  s. (c) Tensile stress distribution on the  
 513 impregnated fiber tow at  $t = 0.1$  s. Reprinted with permission from reference [76], Copyright  
 514 2017, MDPI.

515  
 516 The above studies model the fiber tow as a rigid sliding surface, or a thread constituted by  
 517 rigidly connected spheres. However, this assumption is not valid for the in-situ impregnation-  
 518 based CFRP-AM process, in which polymer melts penetrate into dry fiber bundles during  
 519 the printing process. The degree of impregnation greatly affects the mechanical properties  
 520 of continuous fiber-reinforced composites fabricated by FDM. Although related modeling  
 521 methods have not been reported yet, the existing method for the conventional fiber  
 522 impregnation process can be referred to. In general, there are two methods to model the  
 523 fiber impregnation process. One is to treat the fiber tow as a moving porous medium [82–  
 524 84]. This method ignores the local details of the fiber bundle and defines an anisotropic  
 525 resistance to the fiber region. Another method considers fibers with a microscopic 2D model  
 526 [85], as shown in Fig. 12, where fibers are assumed as rigid surfaces in a micro-scale CFD  
 527 simulation, and it is combined with macro-scale CFD simulation to investigate the fiber  
 528 impregnation process. Meanwhile, the pressure and velocity distribution from the macro-  
 529 scale model are coupled with the micro-scale model. Consequently, the fiber impregnation  
 530 in the FDM process of dry fiber co-extrusion can be thoroughly studied through the methods  
 531 above [82–85]. However, these methods [82–85] fail to consider the impact of fiber  
 532 deformation under the force from the fluid. Thus, new modeling methods are desired to be  
 533 developed to explain the phenomenon of fiber deformation and simulate the complete  
 534 printing process for CRRP.



535

536 Fig. 12. Frame diagram of numerical simulation for fiber impregnation. Reprinted with  
 537 permission from reference [85], Copyright 2021, Elsevier.

## 538 5. Conclusion

539 This study provides a comprehensive review of recent advances in melt flow simulations of  
 540 the FDM process for polymers and their composites. Analytical and numerical methods for  
 541 modeling the extrusion and deposition processes are summarized. Also, it outlines and  
 542 discusses the process simulation for fiber composites which is an emerging research field.  
 543 Finally, the following challenges are identified for inspiring future research:

544 a) Although various non-isothermal, nonlinear material properties have been thoroughly  
 545 considered in current research, further research is needed to deeply understand the effects  
 546 of polymer viscoelasticity on FDM process simulation and account for phenomena such as  
 547 retracting and restarting.

548 b) Recent evidence suggests that the gap between the filament and the hot-end makes  
 549 considerable impacts in the extrusion process of FDM. However, simulations and  
 550 experiments still differ in their results due to the omission of the curvature of the filament.

551 c) For the FDM numerical simulation research of fiber-reinforced composites, in order to  
 552 simulate the whole deposition process and spatial distribution of fiber orientation, it is urgent  
 553 to transition from a 2D model to a 3D model, in which characterizing the fiber interaction  
 554 and decreasing the computation cost are possible directions.

555 d) In the perspective of FDM simulation for continuous fiber-reinforced composites, there  
 556 is a vast and undeniable research gap in modeling the fiber and its interaction with resin in  
 557 the flow field. This includes studying factors such as fiber deformation and fiber  
 558 impregnation.

559 e) There is a need to explore additional applications of simulation models. For instance,  
 560 these simulations can play a critical role in the optimization of the print head structure and  
 561 printing process parameters for emerging FDM processes, e.g., for fiber-reinforced  
 562 composites.

563 **Nomenclature**

564 Stated material properties and printing parameters:

$\mu$	Dynamic viscosity
$\dot{\gamma}$	Shear rate
$\rho$	Density
$C_p$	Specific heat capacity
$h$	Melting enthalpy
$\kappa$	Thermal conductivity
$T_{glass}$	Glass transition temperature
$\dot{q}_{rad}$	Thermal radiation intensity
$g$	Gap height (nozzle position above the substrate)
$D$	Nozzle diameter
$g/D$	Normalized gap height
$u$	Feed rate (the velocity of filament entering the hot-end)
$U$	Extrusion volumetric flux (average velocity inside the nozzle)
$V$	Printing head velocity (printing velocity)
$V/U$	Normalized printing head velocity
$T$	Temperature
$p$	Pressure

565

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570

## 571 **Appendix**

### 572 **Governing equations**

573 The general form of governing equations for non-isothermal, viscous, incompressible, non-  
574 Newtonian fluid flow are as follows:

$$575 \quad \nabla \cdot \mathbf{u} = 0 \quad (1.)$$

$$576 \quad \rho \frac{D\mathbf{u}}{Dt} = -\nabla p + \nabla \cdot \bar{\boldsymbol{\tau}} + \rho \mathbf{g} \quad (2.)$$

$$577 \quad \rho C_p \frac{DT}{Dt} = k \nabla^2 T + \bar{\boldsymbol{\tau}} : \nabla \mathbf{u} \quad (3.)$$

578 Where  $\mathbf{u}$  is the velocity vector;  $\rho$  is the fluid density;  $t$  is the time;  $p$  is the pressure;  $\bar{\boldsymbol{\tau}}$  is the  
579 constitutive stress tensor;  $\mathbf{g}$  is the gravitational acceleration vector;  $C_p$  is the specific heat;  
580  $T$  is the temperature; and  $k$  is the thermal conductivity of the polymer.

581

### 582 **Constituted model**

#### 583 **Generalized Newtonian Fluid (GNF) models mentioned above:**

584 The constitutive stress tensor  $\bar{\boldsymbol{\tau}}$  in Eqs. (1)-(3) typically depends on the instantaneous strain  
585 rate for GNF models, expressed as:

$$586 \quad \bar{\boldsymbol{\tau}} = \eta \mathbf{D} \quad (4.)$$

587 Where  $\mathbf{D}$  is the instantaneous strain rate, calculated as:  $\mathbf{D} = \frac{1}{2}(\nabla \mathbf{u} + \nabla \mathbf{u}^T)$ , and  $\eta$  is the  
588 shear viscosity, described by GNF models as follows:

#### 589 **Power-Law model**

$$590 \quad \eta(\dot{\gamma}) = k(\dot{\gamma})^{n-1} \quad (5.)$$

591 Where  $\eta$  is the shear viscosity,  $\dot{\gamma}$  is the true shear rate,  $k$  is the material consistency, and  $n$   
592 is the power law index.

#### 593 **Carreau-Yasuda model**

$$594 \quad \eta(\dot{\gamma}) = \eta_{\infty} + (\eta_0 - \eta_{\infty}) [1 + (\lambda \dot{\gamma})^a]^{n-\frac{1}{a}} \quad (6.)$$

595 Where  $\eta(\dot{\gamma})$ ,  $\eta_{\infty}$ ,  $\eta_0$ ,  $\lambda$ ,  $a$ , and  $n$  are the shear viscosity, infinite-shear viscosity, zero-shear  
596 viscosity, relaxation time, transition index, and power-law index.

#### 597 **Cross model**

$$598 \quad \eta(\dot{\gamma}) = \frac{\eta_0}{[1 + ((\eta_0/\tau^*)\dot{\gamma})^{1-n}]} \quad (7.)$$

599 Here  $\eta(\dot{\gamma})$ ,  $\eta_0$ , are the shear viscosity, and zero-shear viscosity respectively, and  $\tau^*$  is the  
600 critical shear stress at the transition from the Newtonian plateau,  $n$  is the power-law index.

601

#### 602 **Viscoelastic models mentioned above:**

603 For the viscoelastic constitutive model, accounting for memory effects (indicating internal  
 604 stress of the fluid depends not only on the current deformation rate but also on its past  
 605 deformation history), the constitutive stress tensor  $\tau_s$  in Eqs. (1)-(3) is typically decomposed  
 606 into two parts as:

$$607 \quad \bar{\tau} = \tau_s + \tau_p \quad (8.)$$

608 where  $\tau_s$  is commonly called the solvent stress contribution, and  $\tau_p$  is called the polymer  
 609 stress contribution. The solvent stress contribution depends on the instantaneous strain rate  
 610  $\mathbf{D}$  and is general modelled with newton's law, calculated as:  $\tau_s = \eta_s \mathbf{D}$ ; where  $\eta_s$  is the  
 611 solvent viscosity contribution. And  $\tau_p$  the polymer stress contribution is modeled with  
 612 viscoelastic models as follows:

613 **Phan-Thien-Tanner (PTT) model** [86,87]

$$614 \quad \exp\left[\frac{\varepsilon\lambda}{\eta_p} \text{tr}(\tau_p)\right] \tau_p + \lambda \left[ \left(1 - \frac{\xi}{2}\right) \tau_p + \frac{\xi}{2} \tau_p \right] = 2\eta_p \mathbf{D} \quad (9.)$$

615 Where,  $\xi$  and  $\varepsilon$  are material parameters that control, respectively, the shear viscosity and  
 616 elongational behavior, and  $\lambda$ ,  $\eta_p$  are the relaxation time, and polymer viscosity contribution.

617 **Giesekus model**

$$618 \quad \tau_p + \lambda \tau_{pk}^{\nabla} + \alpha \frac{\lambda}{\eta_p} \tau_p \cdot \tau_p^T = 2\eta_p \mathbf{D} \quad (10.)$$

619 where  $\lambda$ ,  $\alpha$ ,  $\eta_p$  are the relaxation time, mobility factor, and polymer viscosity contribution,  
 620 respectively.

621 **K-BKZ model** [88]

$$622 \quad \tau_p = \frac{1}{1-\theta} \int_{-\infty}^t \sum_{k=1}^N \frac{a_k}{\lambda_k} \exp\left(-\frac{t-t'}{\lambda_k}\right) \left( \frac{\alpha}{(\alpha-3) + \beta I_{C^{-1}} + (1-\beta) I_C} \right) [C_t^{-1}(t') + \theta C_t(t')] dt' \quad (11.)$$

623 where  $a_k$  and  $\lambda_k$  are the relaxation modulus and relaxation time for mode k, N is the number  
 624 of relaxation modes, t is the current time.  $\alpha$  and  $\beta$  are non-linear material constants,  $\theta$  is a  
 625 scalar parameter that controls the ratio of the normal stress differences,  $I_C$  and  $I_{C^{-1}}$  are the  
 626 first invariants of the Cauchy-Green strain tensor  $C_t$  and its inverse  $C_t^{-1}$ , the Finger strain  
 627 tensor.

628

629 **Time-Temperature Equivalence Principle mentioned above:**

630 The temperature-dependent viscosity is often modeled as:

$$631 \quad \eta(\dot{\gamma}, T) = H(T) \eta(\dot{\gamma}) \quad (12.)$$

632

633 Where  $H(T)$  is the temperature shift factor that modifies the viscosity according to the  
 634 polymer temperature, following the time-temperature equivalence principle as follows:

635 **Arrhenius Equation**

$$636 \quad H(T) = \exp\left[\frac{E_a}{R} \left(\frac{1}{T} - \frac{1}{T_{ref}}\right)\right] \quad (13.)$$

637 where  $T_{ref}$  is the reference temperature,  $E_a$  is the activation energy, and R is the gas



638 constant.

639 **William-Landel-Ferry (WLF) Equation**

640 
$$H(T) = \exp \left[ - \frac{C_1(T - T_{ref})}{C_2 + (T - T_{ref})} \right] \quad (14.)$$

641

642 where  $C_1$  and  $C_2$  are the fitting constants,  $T_{ref}$  is reference temperature.

643

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